

A Dimensionless Gravitational State Parameter from a Planck-Scale $U(1)$ Lattice

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We derive a dimensionless gravitational state parameter R from a discrete $U(1)$ lattice with Planck spacing. From $H = J \sum [1 - \cos(\Delta\varphi)]$ with $J = \hbar c / \ell_P$, point defects interact via the 3D Coulomb potential, yielding Newton's law with $n = m/m_P$. The parameter $R \equiv F r^2 / F_P$ ($F_P = c^4 / G$) decomposes into a defect rigidity term $R_m = n^2 r / \ell_P$ and the lowest-order compact $U(1)$ vacuum contribution $2\pi \ell_P / r$:

$$R = \left(\frac{m}{m_P} \right)^2 \frac{r}{\ell_P} - 2\pi \frac{\ell_P}{r}. \quad (1)$$

$R > 0$: classical attraction. $R < 0$: repulsion. $R = 0$ defines a critical line $mr = \sqrt{2\pi} \hbar / c$. The effective force $F = F_P R / r^2$ acquires a mass-independent r^{-4} correction. At laboratory scales this effective correction is estimated to be $\sim 10^{-130}$ N.

I. INTRODUCTION

Quantum gravity lacks a dimensionless parameter characterizing the net gravitational state of a system. Here we derive such a parameter from a discrete $U(1)$ lattice [3, 4]. The parameter R emerges from the competition between defect phase rigidity and compact phase-cycle response.

II. LATTICE FOUNDATION

Space is a 3D grid of spacing $\ell_P = \sqrt{\hbar G / c^3}$ with $U(1)$ phases and Hamiltonian

$$H = J \sum_{\langle i,j \rangle} [1 - \cos(\varphi_i - \varphi_j)], \quad J = \frac{\hbar c}{\ell_P}. \quad (2)$$

A defect satisfies $\oint d\varphi = 2\pi n$. In the continuum limit, the field equation is

$$\nabla^2 \varphi = -4\pi n \delta^{(3)}(r), \quad \varphi(r) = \frac{n}{r}. \quad (3)$$

The interaction potential $V = J n_1 n_2 / r$ follows from the overlap of the two defect phase fields. The force magnitude is

$$|F_{\text{Newton}}| = \left| \frac{dV}{dr} \right| = \hbar c \frac{n_1 n_2}{r^2}. \quad (4)$$

Matching to $F = G m_1 m_2 / r^2$ fixes

$$\boxed{n = \frac{m}{m_P}}, \quad m_P \equiv \sqrt{\frac{\hbar c}{G}}. \quad (5)$$

III. DEFINITION OF R

Define the Planck force $F_P \equiv J / \ell_P = c^4 / G$. The normalized gravitational state parameter is

$$\boxed{R \equiv \frac{F r^2}{F_P}}. \quad (6)$$

This definition removes the geometric $1/r^2$ attenuation common to all 3D point-source forces, exposing the local scale-dependent physics encoded in R . The force is recovered as $F = F_P R / r^2$.

IV. DECOMPOSITION OF R

A. Effective decomposition

The parameter R decomposes into a defect rigidity term and a vacuum response term:

$$\boxed{R = R_m - R_0}, \quad (7)$$

where R_m and R_0 are specified below. The additive structure is the simplest ansatz yielding a well-defined $R = 0$ crossover.

B. Defect term R_m

A defect of winding n produces $|\nabla \varphi| \sim n/r$. The elastic energy $E \sim (J/2) \int (\nabla \varphi)^2 d^3x \sim J n^2 r / \ell_P$. Dividing by $F_P \ell_P = J$ gives the dimensionless contribution

$$\boxed{R_m(r) = n^2 \frac{r}{\ell_P} = \frac{m^2}{m_P^2} \frac{r}{\ell_P}}. \quad (8)$$

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C. Vacuum term R_0

The vacuum's coarse-grained phase response at scale r is governed by the dominant lattice mode $k \sim 1/r$, frequency $\omega = ck \sim c/r$. Over one Planck time $t_P = \ell_P/c$, the accumulated phase advance is $\omega t_P = \ell_P/r$. The lowest-order compact $U(1)$ contribution consistent with the lattice scale is

$$R_0(r) = 2\pi \frac{\ell_P}{r}. \quad (9)$$

The 2π normalization follows from choosing the standard $U(1)$ phase period 2π .

D. Explicit form

$$R(m, r) = \frac{m^2}{m_P^2} \frac{r}{\ell_P} - 2\pi \frac{\ell_P}{r}. \quad (10)$$

V. PROPERTIES

A. Sign regimes

$R > 0$: defect rigidity dominates; the net interaction is attractive. $R < 0$: vacuum response dominates; the interaction is repulsive.

B. Critical condition

$$R = 0 \iff mr = \sqrt{2\pi} \frac{\hbar}{c} \approx 8.82 \times 10^{-43} \text{ kg} \cdot \text{m}. \quad (11)$$

For $m \gtrsim 10^{-8}$ kg, the critical radius $r = \sqrt{2\pi}\hbar/(mc)$ falls below ℓ_P ; such masses are in the $R > 0$ regime at all physically accessible scales.

C. Asymptotics

Large r : $R \approx (m/m_P)^2(r/\ell_P) \gg 0$, recovering the Newtonian regime. Planck scale: $R \sim (m/m_P)^2 - 2\pi$, where the vacuum term is significant.

VI. EFFECTIVE FORCE

From $F = F_P R/r^2$ and Eq. (10):

$$F(r) = G \frac{m_1 m_2}{r^2} - 2\pi \frac{c^4}{G} \left(\frac{\ell_P}{r} \right)^4. \quad (12)$$

The second term is mass-independent, interpreted as an effective vacuum response contribution. At laboratory scales the correction is $\sim 10^{-130}$ N, beyond current sensitivity.

VII. DISCUSSION

We have derived a dimensionless gravitational state parameter R from a discrete $U(1)$ lattice. R decomposes into defect rigidity and compact phase-cycle responses, characterizes the net gravitational state via its sign, and defines a critical line $mr = \sqrt{2\pi}\hbar/c$. The decomposition $R = R_m - R_0$ is an effective-model assumption; the coarse-grained vacuum parametrization $R_0 = 2\pi\ell_P/r$ identifies the lowest-order compact $U(1)$ contribution. The parameter R and the critical condition Eq. (11) provide falsifiable structure for discrete-spacetime phenomenology.

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Appendix A: Additivity and scaling

The definition $R \equiv Fr^2/F_P$ separates the geometric $1/r^2$ from scale-dependent physics. Additivity $R = R_m - R_0$ is the minimal structure producing a well-defined crossover. $R_0 \propto \ell_P/r$ follows from the dominant $k \sim 1/r$ mode; subleading terms $\propto (\ell_P/r)^2$ are Planck-suppressed. The 2π coefficient follows from the chosen compact $U(1)$ phase convention.

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