

Does the blackbody radiation spectrum suggest an intrinsic structure of photons?

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Photons are considered to be elementary bosons in the Standard Model. An assumption that photons are not elementary particles is assessed from an outlook of equilibrium statistical mechanics with insights from computer simulation.

One of the great accomplishments in the last quarter of the nineteenth century was the understanding of how the irreversible macroscopic behavior arises from the time-reversible laws of microscopic physics by Maxwell, Lord Kelvin and Boltzmann. The resolution of apparent paradoxes in statistical mechanics remains controversial even today [1]. Difficulties in our comprehension of how the logic of micro level could be related to the macroscopic behavior are not limited to the origins of the arrow of time. This is where computer simulation of a system consisting of multiple identical elements could be helpful.

The second law of thermodynamics has been illustrated with reversible lattice gas cellular automata (CA). The substitution of individual atoms or molecules with one bit objects on a discrete grid is simplistic but the ability to provide a detailed description of a system – to keep track of every element and actually reverse the evolution of the entire system – is impressive. An expansion of this approach can contribute to understanding of more sophisticated statistical models [2].

It is widely believed that the time reversal invariance plays essentially the same role in classical and quantum theories. The investigation of this paper has started from an attempt to build a reversible working model from the bottom up that would exhibit statistical behavior of massless bosons and generate the blackbody radiation (BBR) spectrum. However, the phase space of photons in Bose-Einstein statistics is not compatible with bijective mapping required for reversible time evolution. This was an obstacle that motivated the search for an alternative solution.

Another reason to search for an alternative to Bose-Einstein statistics came from ontology. According to Leibniz's Principle of the Identity of Indiscernibles (PII), there cannot be separate objects with all the same properties. In physics, classical particles are impenetrable and can be identified by trajectory. Fermions obey Pauli's Exclusion Principle and cannot have all their quantum numbers in common. They are "*weakly discernible*" and the PII still can apply [3, 4]. On the contrary, the elementary bosons can have completely the same sets of properties and be impossible to tell apart and defy the PII. Individual photons are observed in experiments routinely. Are they a good counterexample to the PII as elementary bosons? Should Leibniz's principle be abandoned or, maybe, Bose-Einstein statistics be reevaluated?

I. HISTORICAL ANNOTATIONS

In the derivation of his formula, Planck utilized the product of two factors: the spatial density of radiation energy (in parentheses) and the mean energy U_ν for “monochromatically vibrating resonators” of frequency ν [1]:

$$u_\nu d\nu = \left(\frac{8\pi\nu^2}{c^3} \right) U_\nu d\nu. \quad (1)$$

The resonators can accommodate an integer number of “energy elements” $h\nu$, so the mean energy at temperature T is

$$U_\nu = \frac{h\nu}{e^{h\nu/k_B T} - 1}. \quad (2)$$

The constants above: Boltzmann’s k_B , Planck’s h , and the speed of light c . While the quantization ideas had flourished in a variety of physical applications, indistinguishability of quanta as elements of energy brought questions [6], and the radiation density (number of resonators for each “spectral range”) had not been understood in the same statistical terms as the mean energy.

It took more than 20 years before Bose invented new statistics with “different species of quanta each characterized by the number N_s and energy $h\nu_s$ ($s = 0$ to $s = \infty$)” [7]. He associated the quanta with frequency ν_s with “a cylindrical surface” in a phase space and divided “the total phase space volume into cells of magnitude h^3 ” to get the number of cells for each frequency interval $d\nu$. Bose had arrived at the radiation density in (1) after multiplying the number of cells by the factor 2 to take into account the polarization. The Bose’s derivation of Planck’s formula was “obscure” (in Einstein’s words) and Bose himself was not fully aware of his departure from classical statistics¹.

In 1926 Dirac had incorporated Bose-Einstein statistics into quantum mechanics and linked it with symmetric eigenfunctions [9]. Fowler offered a general form of statistical mechanics in which classical, Bose-Einstein and Fermi-Dirac statistics are special cases [10]. The term *photon* was coined the same year [11] and quickly became popular.

Pauli’s Exclusion Principle was formulated for electrons in 1925, and in 1927 Weyl recognized a connection between Pauli’s principle and the PII. Pauli rejected Weyl’s idea later: “This sounds like a *philosophical* principle and then, it seems to me, there are only two possibilities: a) as such it is wrong; b) it is correct, but *nothing* follows from it for physics ...

¹ “I was not a statistician to the extent of really knowing that I was doing something which was really different from what Boltzmann would have done, from Boltzmann statistics.” (as quoted in [8])

This would really be a strange principle in the philosophy of Leibniz, *which does not hold for all objects* (e.g. not for photons, as Weyl explicitly states) but only for *some*” (as quoted in [12]). The history of quantum statistics and related philosophical questions are presented in [13] and [14].

The first composite photon theory was proposed by de Broglie in 1932 to reconcile photons with Maxwell’s electrodynamics and had nothing to do with statistics. In that search for continuity, the photon consisted of two, then hypothetical, corpuscles: a neutrino and anti-neutrino. Initially, the idea had been pursued by several researchers, but it had not found much traction afterward (see [15] and references therein). In quantum electrodynamics, which is now integrated into a more comprehensive theoretical set of the Standard Model, the photon is still elementary with no known persistent constituents². However, de Broglie's idea may be reintroduced to reconsider the role that gauge bosons are playing within the framework of a quantum field theory [17].

Quantum mechanics was extremely successful in calculating probabilities, but does it point to deterministic or random underlying processes? The Schrödinger equation describes the wave function evolution as deterministic and reversible. Nonetheless, the standard Copenhagen interpretation of quantum mechanics adopted a probabilistic explanation of the wave function (the Born rule) and its indeterministic and irreversible collapse on measurement. Lacking rationale for probabilistic rules was profoundly unsatisfying for some physicists and different interpretations of quantum mechanics were introduced over time (from de Broglie’s pilot wave theory in 1927 to the modern versions of superdeterminism [18], etc.). The quest for understanding quantum foundations is not over.

With the quick development of quantum theories, the physics community became more acceptive of new probabilistic/statistical ideas and new suggestions could be postulated and used without strict causal vetting. When probabilities are primarily measured in experiments, it could be too difficult to recognize what is actually happening in the quantum world and heuristic rules can be adopted³. “Shut up and calculate” [20] could be a justifiable approach for the time being, but it makes sense to revisit older concepts sooner or later and try to understand what is beneath them. Computer simulation could be helpful in determining how the underlying mechanism could work and where those probabilities might be coming from.

² Due to the uncertainty principle, any elementary particle in the quantum field theory, including a photon, can fluctuate into a variety of short-lived virtual states. If a virtual particle interacts with another object, it could expose the structure. The existence of such structure has been well established for photons experimentally at high energies [16].

³ For example, in quantum electrodynamics: “... the price of this great advancement of science is a retreat by physics to the position of being able to calculate only the *probability* that a photon will hit a detector, without offering a good model of how it actually happens ... theoretical physics has given up on that” [19].

II. SIMULATION OF BBR

Massive particles of classical ideal gas exchange energy in reversible elastic collisions. These collisions alone can bring an isolated system to equilibrium. On the other hand, photons do not interact with each other under normal conditions. The mechanism of establishing equilibrium for BBR is emission and absorption of photons in cavity walls (equilibrium with matter). The walls could be regarded as a heat bath for radiation. The number of photons in cavity is not conserved. The preferred framework for the photon gas to explain BBR is Bose-Einstein phase space, and thermalization redistributes photons between the phase space cells. Could this process be reversible like the collisions in an ideal gas and be in agreement with Boltzmann's ideas [1]?

A. The CA model where particles are photons

The integer lattice gas automaton utilized in this investigation (see details in [2]) is based on ideas of continuation of motion (a particle moves in the same direction until it experiences a collision) and detailed balancing (in head-on collision particles are deflected perpendicular). The output of such simulation is defined by initial conditions. Three parameters characterize the system as a whole and are coming from initialization: zero point energy z , the mean occupation number \bar{n} ($\bar{n} > z$) and the step in occupation numbers s . If the lattice sites are initialized with “elements of disorder”, the evolution leads to statistical equilibrium and the most probable exponential distribution is expected for the integer characteristics of the lattice sites.

In this model, structureless gas particles can be seen as photons of the same frequency (a single species of quanta). They are massless and have an assigned fixed energy. For a two-dimensional rectangular CA lattice, each lattice site contains four integers. Each integer accommodates a number of photons moving in one of four possible directions and can be identified with the resonator and/or phase space cell. The automaton redistributes the photons between the resonators/cells. There is no need here in intermediate emission and absorption or cavity walls to thermalize photons. Such an isolated system is fully reversible and can be brought to equilibrium by itself. The Planck's mean energy factor (2) is applicable to it.

If such a system could be expanded for multiple species of quanta, one could get a BBR spectrum. However, a necessary condition for reversibility is conservation of information. It implies bijective uniform mapping from input lattice sites to the same number of output sites in the CA model. Such symmetrical one-to-one mapping cannot be done between different numbers of phase space cells for two or more species of quanta. It would not be possible to establish bijective mapping through any intermediary either – like the cavity walls in traditional understanding of photon thermalization. Thus, the Bose phase space logic for multiple species of quanta is not reversible and could not assist in building a reversible CA.

Indivisible energy is a property of the particle in the CA model and in order to comply with conservation of energy/momentum in each collision, the same value of energy should be

assigned to all the particles. This is another restriction of the model that makes energy exchange impossible between different species of quanta.

With these two restrictions in mind, one can still extend the reversible integer lattice gas model. The lattice can be expanded from two into three dimensions (with six integer characteristics in each lattice site instead of four) or another set of four integers can be added to each site in two-dimensional lattice. The added integers can be assigned to another “breed” of particles with the same energy to make an exchange with the first set possible. All the lattice sites are still uniformly connected with their neighbors, but the elements of the lattice gas model could get a different interpretation. The lattice site attributes can be treated as components of a composite structure. As an example, the next section describes how Einstein’s theory of specific heat can be imitated with such a CA model.

B. Einstein’s specific heat and Wien’s formula for BBR

Einstein made use of Planck’s discontinuity of energy in his theory of specific heat. In his model of a solid, atoms oscillate independently with the same frequency in the three-dimensional lattice. Each of the three degrees of freedom in the oscillations can be associated with a resonator and corresponding mean energy (2). As a result, at low temperature the heat capacity is decreasing but at high temperature the mean energy per atom is still coming to the classical limit $\bar{\varepsilon} = 3k_B T$ ($k_B T$ per degree of freedom).

The detailed energy distribution for atoms (triplets of resonators) in a solid is interesting in the context of this paper. It can be obtained from the integer lattice gas simulation (see Appendix A in [2]) or by using other methods, and it is shaped as Wien’s distribution⁴. To get the distribution from CA simulation, all the integer characteristics of the three-dimensional lattice sites can be divided into three subsets: $E = \{e_1, e_2, \dots, e_N\}$, $M = \{m_1, m_2, \dots, m_N\}$ and $L = \{l_1, l_2, \dots, l_N\}$ with N integers in each one. By combining integers from the subsets, a new variable, ε_i , can be introduced

$$\varepsilon_i = e_i + m_i + l_i. \quad (3)$$

The energy of the triplets (3) produces the Wien’s distribution. The triplet can be seen as a composite structure with the elements of subsets E, M, L as constituents for the structure⁵. In

⁴ Wien introduced his empirical formula in 1896 to describe the BBR spectrum. With quick experimental progress, a call for reassessment came to Planck in 1900, and he improved Wien’s formula.

⁵ If demonstration of reversibility is not a priority, a simple stochastic technique can be used instead of CA simulation to generate the subsets (see Appendix A). While the integer lattice gas provides a working model of a system and the automaton is a reflection of microdynamics, it requires more computational resources and could have a long *relaxation time*. On the other hand, the Monte Carlo approach is a fast way of getting specific distributions while disregarding the cause. Either way, one can produce elements to build the composite structures like (3).

this interpretation the lattice gas particles are not complete physical objects any more (like the photons as they were interpreted in the previous section) but simple *energy bits* for the constituents. Each integer in (3) can be associated with a translational degree of freedom in Einstein’s model of a solid. Every energy bit has the sole energy ε_z . Table I provides a short summary of the two interpretations of the lattice gas in sections A and B.

TABLE I. Two interpretations of the integer lattice gas

| | A (basic) | B (composite structures) |
|---|--|--|
| particle | elementary photon | energy bit |
| integer characteristic of lattice site | resonator or phase space cell | resonator or a constituent of composite object |
| number of particles in integer characteristic | number of photons associated with the resonator or phase space cell | energy of the resonator or constituent |
| good for | single mode of radiation; most probable distributions | detailed look at Einstein’s specific heat; Wien’s distribution |
| what is questionable | cannot be used with Bose-Einstein phase space; not consistent with the PII | |

C. Photon structure

The triplets (3) can generate Wien’s distribution. Could a similar structure bring the energy distribution closer to the Planckian spectrum with the corresponding average energy per photon, $\bar{\varepsilon} \approx 2.7k_bT$? It is only about 10% less than for the triplet. A duplet of noninteracting constituents would not have sufficient average energy for BBR.

By examining statistics for different combinations of two numbers – one from each of the subsets E and M – it has been found in this study that the sum of both plus the geometric mean,

$$\varepsilon_i = e_i + m_i + \sqrt{e_i m_i}, \quad (4)$$

produces a distribution that is close in shape to Planck’s law. It is presented in Fig. 1. Planck’s law function graph is for the fixed number of photons (Appendix B) to make a

proper comparison to the energy spectrum of the same number of *constructed* photons (4). The temperature is the same in simulated exponential distributions for the subsets E and M and in function (B5). The units of energy ($\varepsilon_z = 1$) are the same for both.

The energy quantization for the constituents would cause degradation of the energy distribution for composite photons at low mean occupation numbers that would become pronounced at $\bar{n} < 30$. Such degradation is not foreseen in the standard BBR explanation. To avoid this effect, the presented computations are for $\bar{n} \approx k_B T / \varepsilon_z = 600$. Other simulation parameters are $z = 1$ and $s = 1$. It is assumed: each constituent holds at least one energy bit (zero-point energy is equal to unity – no emptiness).

The two independent variables in (4) can be understood as two constituents of a photon for which energy can vary. The geometric mean can be seen as interaction energy between the two. It is fully defined by energy of two constituents and is not coming as another degree of freedom. From the perspective of the structure, the average energy per photon in the BBR spectrum, $\bar{\varepsilon} \approx 2.7k_B T$, comprises the average energy of two constituents, $2k_B T$, and the interaction, $\sim 0.7k_B T$. If energy is defined as a total number of particles in the CA system, like it was done before this section, the conservation of energy is not in question. On the other hand, the total interaction energy for the system in this section, $\sum_N \sqrt{e_i m_i}$, will fluctuate with evolution steps.

The CA rule updates all cells in one step. It can be seen as working in the discrete time domain where the logic of events is not spatial. The CA grid is organized into a number of dimensions but does not provide continuous symmetries or other features one would expect from a fully fledged space or spacetime. However, the physical phenomenon of BBR and Planck's law define not only the shape of energy distribution but the number of photons in the unit volume for the temperature (B4). It can be used to introduce distance and volume as secondary attributes into the system⁶.

⁶ The system can be “*spatially extended*” [21] by assigning length to connections to neighboring cells in the CA grid. According to (B4), one photon on average takes a volume of a cubic box with a side length proportional to $\frac{hc}{k_B T}$. The length (distance between cells) can be attributed to a photon in this model and defined as the

function – $d = \frac{w}{\varepsilon_i}$, where w is a constant – to satisfy spatial requirements of BBR. This function is analogous

to the photon's wavelength-energy relationship in conventional terms: $\lambda = \frac{hc}{\varepsilon}$. The spatial locations of cells in

such a system are relative and would fluctuate with changes in energy of each cell with evolution steps. The assumption of nonzero positive minimal energy for photons is required for this distance function. Physical space in this model is curved by radiation like it is curved by masses in General Relativity and there is no “empty” space. A similar assumption of positive minimal energy was made for the constituents to form energy distribution in Fig. 1. The CA models are considered as candidates for emergent space of fundamental physics [22, 18].

The two constituents, with a quite arbitrarily introduced interaction, bring the energy distribution close to Planck's radiation law over a broad spectral range. This interaction could point to a force between the constituents but it is not originated in the generic microdynamics used to produce the constituents. This structure is reminiscent of mesons in the Standard Model: one quark and one antiquark bound together by the strong interaction. De Broglie's attempt to reconcile photons with Maxwell's electrodynamics brought him to the first composite photon theory. This paper represents another search for continuity, now in statistics, but again, points to the structure that is akin to de Broglie's.

III. CONCLUSION

Bose-Einstein statistics was first introduced specifically to explain Planck's law. Further advances in particle physics discovered multiple other bosons in addition to photons. The majority of all bosons are believed to be constructed from an even number of quarks/antiquarks or other particles, while photons, along with other gauge bosons and Higgs bosons, are still regarded as elementary particles in the Standard Model.

The Bose phase space is not compatible with the time reversal invariance of microscopic processes in classical and quantum mechanics and cannot be used in the bottom-up approach to photon statistics. A possible alternative justification of the Planckian spectrum can come from the intrinsic structure. If one accepts BBR spectrum as a manifestation of the photon's structure, one might also infer that other fundamental bosons have a similar intrinsic organization and are not elementary.

If all bosons were not elementary, Leibniz's Principle of the Identity of Indiscernibles would be vindicated in particle physics, since "the only cases in which the status of quantum particles as objects is seriously in question are ... *elementary* bosons – bosons (supposedly) with no internal fermionic structure" [3].

The understanding of quantum statistics' origins plays a key role in constructing fundamental unified theories beyond the Standard Model. In supersymmetry, bosons and fermions are treated as fundamental particles. Each fundamental particle from one group has an associated so-called "superpartner" in another group. Supersymmetry predicts a large number of undiscovered elementary particles. With no fundamental bosons, there would not be a rationale for this kind of symmetry. Some alternative theories (like Spinor Gravity [23] and causal fermion systems [24]) treat the fermions as fundamental particles and consider the bosons as composite objects made of an even number of fermions.

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APPENDIX A: STOCHASTIC SIMULATION

A Monte-Carlo technique similar to the one used in numerical integration can be deployed to populate arrays (or sets) of N integers and obtain distribution of a given shape $y = f(x)$ for those numbers.

To do so, pairs of real pseudorandom numbers (R_x and R_y) can be generated as points in the rectangular region that entirely covers the function graph $y = f(x)$. If R_y falls below $f(R_x)$, the value R_x is rounded up to the nearest integer and populates one element in the array of integers. Otherwise, the pair of pseudorandom numbers is discarded. The cycle is repeated till all elements of the array are filled.

Each array would form a most probable (Boltzmann) distribution defined by the exponential function $y = f(x) = e^{-x\varepsilon_z/k_B T}$. The integer would stand for the number of energy bits in the constituent of composite structure. The rounding up makes the spectrum discrete and sets its minimum to unity.

APPENDIX B: PLANCK'S RADIATION LAW FOR A FIXED NUMBER OF PHOTONS

With photon energy ε instead of $h\nu$, Planck's law for a unit volume can be written as:

$$u(T, \varepsilon) = \frac{8\pi}{(hc)^3} \frac{\varepsilon^3}{e^{\varepsilon/k_B T} - 1}. \quad (\text{B1})$$

The corresponding number of photons is distributed with energy as:

$$n(T, \varepsilon) = \frac{u(T, \varepsilon)}{\varepsilon} = \frac{8\pi}{(hc)^3} \frac{\varepsilon^2}{e^{\varepsilon/k_B T} - 1}. \quad (\text{B2})$$

And the total number of photons in the unit volume at temperature T can be found from integration:

$$N_0 = \int_0^{\infty} n(T, \varepsilon) d\varepsilon = \frac{8\pi}{(hc)^3} \int_0^{\infty} \frac{\varepsilon^2 d\varepsilon}{e^{\varepsilon/k_B T} - 1}.$$

Let $x = \varepsilon/k_B T$ and take into account that

$$\int_0^{\infty} \frac{x^2 dx}{e^x - 1} = 2\zeta(3) \approx 2.404, \quad (\text{B3})$$

where $\zeta(3)$ is the Riemann zeta function, also known as Apéry's constant. (The stochastic procedure like the one in Appendix A could be utilized to compute this integral as well.) This results in the spatial factor in Planck's law,

$$N_0 \approx 2.404 \frac{8\pi(k_B T)^3}{(hc)^3}, \quad (\text{B4})$$

with the fixed number of photons, N , distributed by energy as:

$$f(T, \varepsilon) = \frac{N}{N_0} n(T, \varepsilon) \approx \frac{N}{2.404(k_B T)^3} \frac{\varepsilon^2}{e^{\varepsilon/k_B T} - 1}. \quad (\text{B5})$$

- [1] S. Goldstein, *Boltzmann's approach to statistical mechanics*, in: *Chance in Physics: Foundations and Perspectives*, edited by J. Bricmont, et al., *Lecture Notes in Physics*, 574, Springer-Verlag, Berlin, 2001, p.39-68; arXiv:cond-mat/0105242.
- [2] A. Khaneles, *Int. J. Appl. Math. Stat. (IJAMAS)*, Vol. 4, **M06**, 44-57 (2006); cond-mat/0512292.
- [3] S. Saunders, *Analysis* **66**, 52-63, (2006); philsci-archive.pitt.edu/2623/
- [4] A. Guay and T. Pradeu, eds., *Individuals Across the Sciences*, (Oxford University Press, 2015).
- [5] M. Planck, *Verhandlungen der Deutschen Physikalischen Gesellschaft*, **2**, 237 (1900) [D. ter Haar, S. G. Brush, *Planck's Original Papers in Quantum Physics*, (Taylor and Francis, London, 1972), p.38].
- [6] S. Saunders, 'Indistinguishability', in *Oxford Handbook of Philosophy of Physics*, R. Batterman (ed.), (Oxford University Press, 2013); <http://philsci-archive.pitt.edu/12448/>
- [7] S.N. Bose, *Z. Phys.* **26**, 178 (1924); [O. Theimer and B. Ram, *Am. J. Phys.*, **44**, 1056 (1976)].
- [8] A. Pais, *Subtle is the Lord...*, (Oxford University Press, New York, 1982), p.424.
- [9] P.A.M. Dirac, *Proc. R. Soc. A*, **112**, 661-677, (1926).
- [10] R.H. Fowler, *Proc. R. Soc. A*, **113**, 432-449, (1926).
- [11] H. Kragh, *Photon: New light on an old name*, arXiv:1401.0293.
- [12] C. P. Enz, *No Time to be Brief: A Scientific Biography of Wolfgang Pauli*. (New York: Oxford University Press, 2002).
- [13] M. Jammer, *The Conceptual Development of Quantum Mechanics*, (New York: McGraw-Hill, 1966).
- [14] S. French and D. Krause, *Identity in Physics: A Historical, Philosophical and Formal Analysis*, (Oxford: Oxford University Press, 2006).
- [15] W.A. Perkins, *Journal of Modern Physics*, **5**, 2089-2105, (2014); physics.gen-ph/1503.00661v1.
- [16] C. Berger, *Journal of Modern Physics*, **6**, 1023-1043, (2015); arXiv:1404.3551.
- [17] H. Stumpf and T. Borne, *Annales de Fondation Louis de Broglie*, **26**, 429-448, (2001).
- [18] G. 't Hooft, *The Cellular Automaton Interpretation of Quantum Mechanics*, (Springer, 2016).
- [19] R.P. Feynman, *QED: The Strange Theory of Light and Matter*, (Princeton University Press, Princeton, 1985), p. 37, 82.
- [20] N.D. Mermin, *Phys. Today*, **42**, 4, 9 (1989).
- [21] A. J. Kox, Martin J. Klein, and Robert Schulmann (eds.), *The Collected Papers of Albert Einstein. Vol. 6. The Berlin Years: Writings, 1914-1917*, (Princeton University Press, Princeton, 1996), Doc.42, p.419.
- [22] I. Licata (ed.), *Beyond peaceful coexistence: the emergence of space, time and quantum*, (Imperial College Press, London, 2016).
- [23] A. Hebecker, C. Wetterich, *Phys. Lett. B*, **574**, 269, (2003); arXiv:hep-th/0307109.
- [24] F. Finster, J. Kleiner, *Causal fermion systems as a candidate for a unified physical theory*, *J. Phys. Conf. Ser.* **626**, 012020 (2015); arXiv:1502.03587.

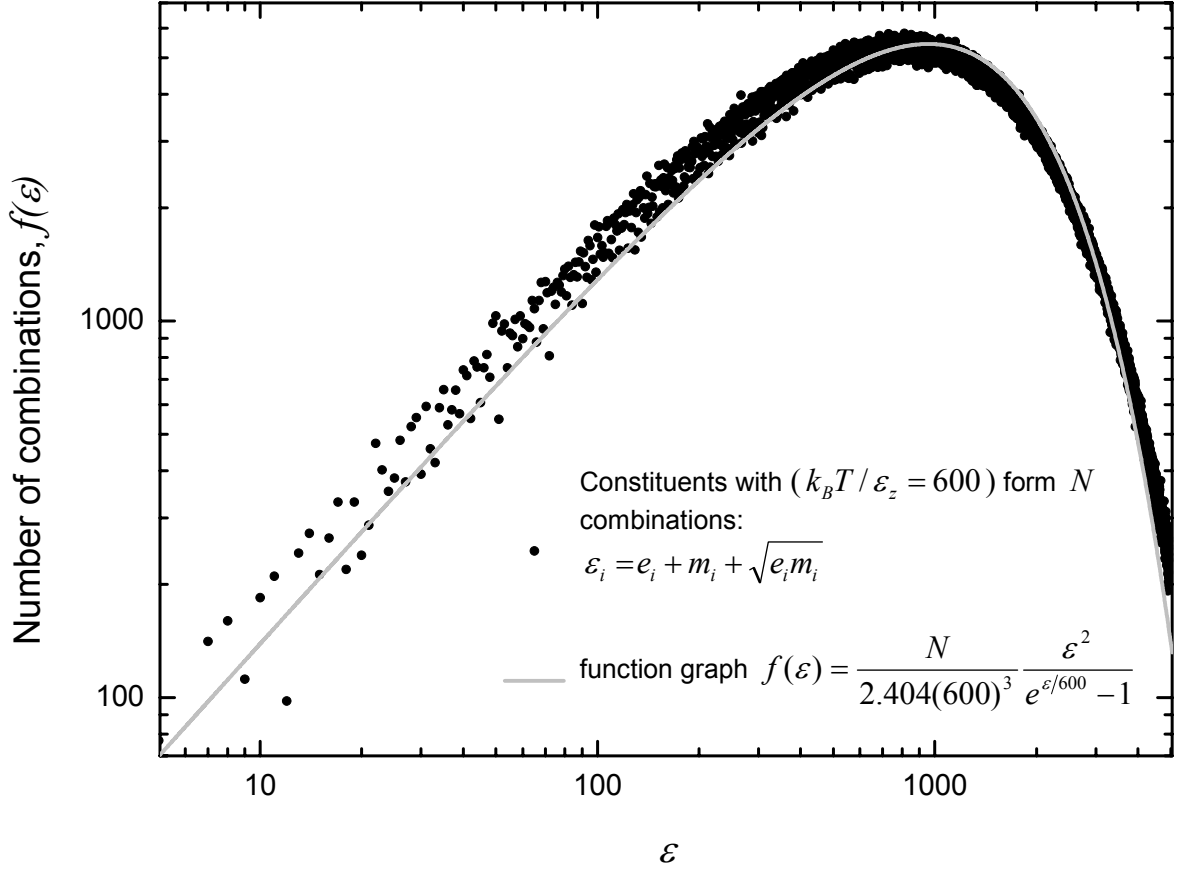


FIG. 1. The number of combinations ($\varepsilon_i = e_i + m_i + \sqrt{e_i m_i}$) in the distribution is N . The function graph is for the same number of photons. The arrays of constituents used in this simulation are composed of $N \sim 12 \cdot 10^6$ integers each. The simulation parameters are: the mean occupation number $\bar{n} \approx k_B T / \varepsilon_z = 600$, zero point energy $z = 1$ and step in occupation numbers $s = 1$.