

# **The Electro-Magnetic Field Equation and Electro-Magnetic Field Transformation in Rindler spacetime**

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## **ABSTRACT**

In the general relativity theory, we find the electro-magnetic field transformation and the electro-magnetic field equation (Maxwell equation) in Rindler spacetime.

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## 1. Introduction

In the general relativity theory, we find the electro-magnetic field transformation and the electro-magnetic field equation (Maxwell equation) in Rindler spacetime.

The Rindler coordinate is

$$\begin{aligned}
 ct &= \left(\frac{c^2}{a_0} + \xi^1\right) \sinh\left(\frac{a_0 \xi^0}{c}\right) \\
 x &= \left(\frac{c^2}{a_0} + \xi^1\right) \cosh\left(\frac{a_0 \xi^0}{c}\right) - \frac{c^2}{a_0}, \quad y = \xi^2, \quad z = \xi^3
 \end{aligned} \tag{1}$$

In this time, the tetrad  $e^a{}_\mu$  is

$$\begin{aligned}
 d\tau^2 &= dt^2 - \frac{1}{c^2} [dx^2 + dy^2 + dz^2] \\
 &= -\frac{1}{c^2} \eta_{ab} \frac{\partial x^a}{\partial \xi^\mu} \frac{\partial x^b}{\partial \xi^\nu} d\xi^\mu d\xi^\nu \\
 &= -\frac{1}{c^2} \eta_{ab} e^a{}_\mu e^b{}_\nu d\xi^\mu d\xi^\nu = -\frac{1}{c^2} g_{\mu\nu} d\xi^\mu d\xi^\nu, \quad e^a{}_\mu = \frac{\partial x^a}{\partial \xi^\mu} \tag{2}
 \end{aligned}$$

$$e^{\alpha}{}_0(\xi^0) = \frac{\partial x^\alpha}{\partial \xi^0} = \left(\cosh\left(\frac{a_0 \xi^0}{c}\right), \sinh\left(\frac{a_0 \xi^0}{c}\right), 0, 0\right) \tag{3}$$

About  $y$ -axis's and  $z$ -axis's orientation

$$e^{\alpha}{}_2(\xi^0) = \frac{\partial x^\alpha}{\partial \xi^2} = (0, 0, 1, 0), \quad e^{\alpha}{}_3(\xi^0) = \frac{\partial x^\alpha}{\partial \xi^3} = (0, 0, 0, 1) \tag{4}$$

The other unit vector  $e^{\alpha}{}_1(\xi^0)$  is

$$e^{\alpha}{}_1(\xi^0) = \frac{\partial x^\alpha}{\partial \xi^1} = \left(\sinh\left(\frac{a_0 \xi^0}{c}\right), \cosh\left(\frac{a_0 \xi^0}{c}\right), 0, 0\right) \tag{5}$$

Therefore,

$$\begin{aligned}
 cdt &= c \cosh\left(\frac{a_0 \xi^0}{c}\right) d\xi^0 \left(1 + \frac{a_0}{c^2} \xi^1\right) + \sinh\left(\frac{a_0 \xi^0}{c}\right) d\xi^1 \\
 dx &= c \sinh\left(\frac{a_0 \xi^0}{c}\right) d\xi^0 \left(1 + \frac{a_0}{c^2} \xi^1\right) + \cosh\left(\frac{a_0 \xi^0}{c}\right) d\xi^1, \quad dy = d\xi^2, \quad dz = d\xi^3 \tag{6}
 \end{aligned}$$

Hence, the transformation of the electro-magnetic potential  $(\phi, \vec{A})$  in inertial frame and the electro-

magnetic potential  $(\phi_{\xi}, \vec{A}_{\xi})$  is

$$\begin{aligned}\phi &= \cosh\left(\frac{a_0 \xi^0}{c}\right) \left(1 + \frac{a_0}{c^2} \xi^1\right) \phi_{\xi} + \sinh\left(\frac{a_0 \xi^0}{c}\right) A_{\xi^1} \\ A_x &= \sinh\left(\frac{a_0 \xi^0}{c}\right) \left(1 + \frac{a_0}{c^2} \xi^1\right) \phi_{\xi} + \cosh\left(\frac{a_0 \xi^0}{c}\right) A_{\xi^1} \\ A_y &= A_{\xi^2}, A_z = A_{\xi^3}\end{aligned}\tag{7}$$

$$e^a{}_{\mu} e^b{}_{\nu} \eta_{ab} = g_{\mu\nu} \rightarrow A^T \eta A = g$$

$$e_a{}^{\mu} e_b{}^{\nu} g_{\mu\nu} = \eta_{ab} \rightarrow (A^T)^{-1} g A^{-1} = (A^T)^{-1} A^T \eta A A^{-1} = \eta$$

$$e^a{}_{\mu} = \eta^{ab} g_{\mu\nu} e_b{}^{\nu} \rightarrow \eta^{-1} (A^T)^{-1} A^T \eta A = A = \eta^{-1} (A^T)^{-1} g$$

$$\begin{aligned}\begin{pmatrix} cdt \\ dx \\ dy \\ dz \end{pmatrix} &= \begin{pmatrix} \cosh\left(\frac{a_0 \xi^0}{c}\right) \left(1 + \frac{a_0 \xi^1}{c^2}\right) & \sinh\left(\frac{a_0 \xi^0}{c}\right) & 0 & 0 \\ \sinh\left(\frac{a_0 \xi^0}{c}\right) \left(1 + \frac{a_0 \xi^1}{c^2}\right) & \cosh\left(\frac{a_0 \xi^0}{c}\right) & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} cd\xi^0 \\ d\xi^1 \\ d\xi^2 \\ dz^3 \end{pmatrix} \\ &= A \begin{pmatrix} cd\xi^0 \\ d\xi^1 \\ d\xi^2 \\ dz^3 \end{pmatrix}\end{aligned}$$

$$\begin{pmatrix} \frac{1}{c} \frac{\partial}{\partial t} \\ \frac{\partial}{\partial x} \\ \frac{\partial}{\partial y} \\ \frac{\partial}{\partial z} \end{pmatrix} = (A^{-1})^T \begin{pmatrix} \frac{1}{c} \frac{\partial}{\partial \xi^0} \\ \frac{\partial}{\partial \xi^1} \\ \frac{\partial}{\partial \xi^2} \\ \frac{\partial}{\partial \xi^3} \end{pmatrix} = (A^T)^{-1} \begin{pmatrix} \frac{1}{c} \frac{\partial}{\partial \xi^0} \\ \frac{\partial}{\partial \xi^1} \\ \frac{\partial}{\partial \xi^2} \\ \frac{\partial}{\partial \xi^3} \end{pmatrix}$$

$$= \begin{pmatrix} \frac{\cosh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} & -\sinh(\frac{a_0 \xi^0}{c}) & 0 & 0 \\ -\frac{\sinh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} & \cosh(\frac{a_0 \xi^0}{c}) & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \frac{1}{c} \frac{\partial}{\partial \xi^0} \\ \frac{\partial}{\partial \xi^1} \\ \frac{\partial}{\partial \xi^2} \\ \frac{\partial}{\partial \xi^3} \end{pmatrix}$$

$$\frac{1}{c} \frac{\partial}{\partial t} = \frac{\partial \xi^0}{\partial t} \frac{1}{c} \frac{\partial}{\partial \xi^0} + \frac{\partial \xi^1}{\partial t} \frac{\partial}{\partial \xi^1}$$

$$= \frac{\cosh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{\partial \xi^0} - \sinh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \xi^1}$$

$$\frac{\partial}{\partial x} = \frac{\partial \xi^0}{\partial x} \frac{1}{c} \frac{\partial}{\partial \xi^0} + \frac{\partial \xi^1}{\partial x} \frac{\partial}{\partial \xi^1}$$

$$= -\frac{\sinh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{\partial \xi^0} + \cosh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \xi^1}$$

$$\frac{\partial}{\partial y} = \frac{\partial}{\partial \xi^2}, \quad \frac{\partial}{\partial z} = \frac{\partial}{\partial \xi^3} \quad (8)$$

## 2. Electro-magnetic Field in the Rindler space-time

The electro-magnetic field  $(\vec{E}, \vec{B})$  is in the inertial frame,

$$\vec{E} = -\vec{\nabla} \phi - \frac{\partial \vec{A}}{\partial t}, \quad \vec{B} = \vec{\nabla} \times \vec{A} \quad (9)$$

We define the electro-magnetic field  $(\vec{E}_\xi, \vec{B}_\xi)$  in Rindler spacetime.

$$\vec{E}_\xi = -\frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \vec{\nabla}_\xi \left\{ \phi_\xi \left( 1 + \frac{a_0 \xi^1}{c^2} \right)^2 \right\} - \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial \vec{A}_\xi}{\partial \xi^0}$$

$$\vec{B}_\xi = \vec{\nabla}_\xi \times \vec{A}_\xi$$

In this time,  $\vec{\nabla}_\xi = (\frac{\partial}{\partial \xi^1}, \frac{\partial}{\partial \xi^2}, \frac{\partial}{\partial \xi^3}), \vec{A}_\xi = (A_{\xi^1}, A_{\xi^2}, A_{\xi^3})$  (10)

Hence, we obtain the transformation of the electro-magnetic field.

$$E_x = -\frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{\partial \xi^1} \{ \phi_\xi (1 + \frac{a_0 \xi^1}{c^2})^2 \} - \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial A_{\xi^1}}{\partial \xi^0} = E_{\xi^1},$$

$$E_y = E_{\xi^2} \cosh(\frac{a_0 \xi^0}{c}) + B_{\xi^3} \sinh(\frac{a_0 \xi^0}{c}),$$

$$E_z = E_{\xi^3} \cosh(\frac{a_0 \xi^0}{c}) - B_{\xi^2} \sinh(\frac{a_0 \xi^0}{c})$$

$$B_x = B_{\xi^1},$$

$$B_y = B_{\xi^2} \cosh(\frac{a_0 \xi^0}{c}) - E_{\xi^3} \sinh(\frac{a_0 \xi^0}{c})$$

$$B_z = B_{\xi^3} \cosh(\frac{a_0 \xi^0}{c}) + E_{\xi^2} \sinh(\frac{a_0 \xi^0}{c}) \quad (11)$$

### 3. Electro-magnetic Field Equation(Maxwell Equation) in the Rindler space-time

Maxwell equation is

$$\vec{\nabla} \cdot \vec{E} = 4\pi\rho \quad (12-i)$$

$$\vec{\nabla} \times \vec{B} = \frac{\partial \vec{E}}{\partial t} + \frac{4\pi}{c} \vec{j} \quad (12-ii)$$

$$\vec{\nabla} \cdot \vec{B} = 0 \quad (12-iii)$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad (12-iv)$$

Therefore, we obtain the electro-magnetic field equation by Eq (12-i), Eq(12-ii), Eq(12-iii), Eq(12-iv) in Rindler spacetime .

$$\vec{\nabla}_\xi \cdot \vec{E}_\xi = 4\pi\rho_\xi \quad (13-i)$$

$$\frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \vec{\nabla}_\xi \times \{ \vec{B}_\xi (1 + \frac{a_0 \xi^1}{c^2}) \} = \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial \vec{E}_\xi}{\partial \xi^0} + \frac{4\pi}{c} \vec{j}_\xi \quad (13-ii)$$

$$\vec{\nabla}_\xi \cdot \vec{B}_\xi = 0 \quad (13-iii)$$

$$\frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \vec{\nabla}_\xi \times \{ \vec{E}_\xi (1 + \frac{a_0 \xi^1}{c^2}) \} = - \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial \vec{B}_\xi}{\partial \xi^0} \quad (13\text{-iv})$$

$$\vec{E}_\xi = (E_{\xi^1}, E_{\xi^2}, E_{\xi^3}), \vec{B}_\xi = (B_{\xi^1}, B_{\xi^2}, B_{\xi^3}), \vec{\nabla}_\xi = (\frac{\partial}{\partial \xi^1}, \frac{\partial}{\partial \xi^2}, \frac{\partial}{\partial \xi^3})$$

Hence,

$$\begin{aligned} \rho &= \rho_\xi \cosh(\frac{a_0 \xi^0}{c}) + \frac{j_{\xi^1}}{c} \sinh(\frac{a_0 \xi^0}{c}) \\ j_x &= j_{\xi^1} \cosh(\frac{a_0 \xi^0}{c}) + c \rho_\xi \sinh(\frac{a_0 \xi^0}{c}), \quad j_y = j_{\xi^2}, j_z = j_{\xi^3} \end{aligned} \quad (14)$$

In this time, the Lorentz gauge is

$$\begin{aligned} \frac{\partial \phi}{\partial t} + \vec{\nabla} \cdot \vec{A} &= \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{\partial \xi^1} \{ A_{\xi^1} (1 + \frac{a_0 \xi^1}{c^2}) \} + \frac{\partial \phi_\xi}{\partial \xi^0} + \frac{\partial A_{\xi^2}}{\partial \xi^2} + \frac{\partial A_{\xi^3}}{\partial \xi^3} \\ &= \frac{\partial \phi_\xi}{\partial \xi^0} + \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \vec{\nabla}_\xi \cdot \{ \vec{A}_\xi (1 + \frac{a_0 \xi^1}{c^2}) \} = 0 \end{aligned} \quad (15)$$

#### 4. Conclusion

In the general relativity theory, we find the electro-magnetic field transformation and the electro-magnetic field equation (Maxwell equation) in Rindler spacetime.

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