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## The special theory of relativity and the law of conservation of momentum

This article attempts to use the law of conservation of momentum of a closed system for determining the values of the constants in the two possible transformations of coordinates and times in inertial reference systems.

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## 1. The introduction

Special theory of relativity can be divided into relativistic kinematics and relativistic dynamics.

Relativistic kinematics, based on the symmetry of space and time [1], [2], [3], [4], [5], the principle of relativity and the principle of invariance of the speed of light [6], [7], [8], [9 ], [10], [11], [12], 13], allows to establish links between the coordinates and times (Lorentz transformation), speeds (the transformation of speeds) and accelerations (the transformation of accelerations) in the inertial reference systems.

The relativistic dynamics, based on the mandatory implementation of the laws of conservation of momentum, angular momentum and energy of a closed system [14], [15], [16], [17], [18], [19], [20], [21], [ 22], [23] in the inertial reference systems, determines the dependences of mass, momentum and energy point of the material body from its speed.

However, in the experiments [24] and the analysis of the results of observations [25], [26] were marked by inconsistency of the actual results the conclusions of the special theory of relativity.

To understand the reason that caused the deviation, can, by analogy with [4] consider the special theory of relativity in a general form with less severe conditions - without the use of the principle of invariance of the speed of light.

## 2. The special theory of relativity in the general form

Suppose that the space is homogeneous and isotropic and time is homogeneous (ie there is a symmetry of space and time).

In considering will only use the principle of relativity, asserts, that in any inertial reference systems all physical phenomena occur equally in the same conditions.

Assume that there are two inertial reference systems, shown in Fig.1, stationary $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and mobile $\mathrm{O}_{2} x_{2} y_{2} z_{2}$, in which:

- similar the axis of the Cartesian coordinate systems $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and $\mathrm{O}_{2} x_{2} y_{2} z_{2}$
are pairs parallel and equally directed;
- system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ moves relative to the system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ with constant speed $V$ along the axis $\mathrm{O}_{1} x_{1}$;
- in both systems as the start timing ( $t_{1}=0$ and $\left.t_{2}=0\right)$ is selected when the origin $\mathrm{O}_{1}$ and $\mathrm{O}_{2}$ of these systems are identical.


Fig. 1
Using the principle of relativity and the symmetry of space and time, by analogy with [27], [8], [9], [10], [18], provides a link between the coordinates $x_{1}$, $y_{1}, z_{1}$ of point A at time $t_{1}$ in a stationary inertial reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and coordinates $x_{2}, y_{2}, z_{2}$ of the same point A in the mobile inertial reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at the time $t_{2}$, corresponding to time $t_{1}$ in the stationary inertial reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ :

$$
\begin{gather*}
x_{1}=\gamma_{V} \cdot\left[x_{2}+\left(V \cdot t_{2}\right)\right]  \tag{1}\\
x_{2}=\gamma_{V} \cdot\left[x_{1}-\left(V \cdot t_{1}\right)\right] \\
y_{1}=y_{2}  \tag{3}\\
z_{1}=z_{2} \tag{4}
\end{gather*}
$$

where: $\gamma_{V}$ - coefficient of proportionality (the transition), which is presumably a function of speed $V$.

From formulas (1) and (2) we can write the dependence for times $t_{1}$ and $t_{2}$ :

$$
\begin{equation*}
t_{1}=\frac{\left(\gamma_{V}^{2}-1\right) \cdot x_{2}}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{2}\right) \tag{5}
\end{equation*}
$$

$$
\begin{equation*}
t_{2}=\frac{\left(1-\gamma_{V}^{2}\right) \cdot x_{1}}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{1}\right) \tag{6}
\end{equation*}
$$

Differentiating equations (1) - (6), we can obtain the relationship between the projections $v_{x 1}$, $v_{y 1}$ and $v_{z 1}$ of the speed of a point on the axis of the Cartesian coordinates in time $t_{1}$ in the stationary inertial reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and similar projections $v_{x 2}, v_{y 2}$ and $v_{z 2}$ of the speed of the same point in the mobile inertial reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{2}$, corresponding to time $t_{1}$ in the stationary inertial reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ :

$$
\begin{align*}
& v_{x 1}=\frac{v_{x 2}+V}{\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{x 2}}{\gamma_{V}^{2} \cdot V}+1}  \tag{7}\\
& v_{x 2}=\frac{v_{x 1}-V}{\frac{\left(1-\gamma_{V}^{2}\right) \cdot v_{x 1}}{\gamma_{V}^{2} \cdot V}+1}  \tag{8}\\
& v_{y 1}=\frac{v_{y 2}}{\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{x 2}}{\gamma_{V} \cdot V}+\gamma_{V}}  \tag{9}\\
& v_{y 2}=\frac{v_{y 1}}{\frac{\left(1-\gamma_{V}^{2}\right) \cdot v_{x 1}}{\gamma_{V} \cdot V}+\gamma_{V}}  \tag{10}\\
& v_{z 1}=\frac{v_{z 2}}{\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{x 2}}{\gamma_{V} \cdot V}+\gamma_{V}}  \tag{11}\\
& v_{z 2}=\frac{v_{z 1}}{\frac{\left(1-\gamma_{V}^{2}\right) \cdot v_{x 1}}{\gamma_{V} \cdot v_{x 1}}+\gamma_{V}} \tag{12}
\end{align*}
$$

Differentiation of equations (7) - (12) and (5) - (6) will be written communication between the projections $a_{x 1}, a_{y 1}$ and $a_{z 1}$ of the acceleration of a point on the axis of the Cartesian coordinates in time $t_{1}$ in the stationary inertial reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and similar projections $a_{x 2}, a_{y 2}$ and $a_{z 2}$ of the acceleration of the same point in the mobile inertial reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{2}$, corresponding to time $t_{1}$ in the stationary inertial reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ :

$$
\begin{gather*}
a_{x 1}=\frac{a_{x 2}}{\gamma_{V}^{3} \cdot\left[\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{x 2}}{\gamma_{V}^{2} \cdot V}+1\right]^{3}}  \tag{13}\\
a_{x 2}=\frac{a_{x 1}}{\gamma_{V}^{3} \cdot\left[\frac{\left(1-\gamma_{V}^{2}\right) \cdot v_{x 1}}{\gamma_{V}^{2} \cdot V}+1\right]^{3}}  \tag{14}\\
a_{y 1}=\frac{\left\{a_{y 2} \cdot\left[\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{x 2}}{\gamma_{V} \cdot V}+\gamma_{V}\right]\right\}-\frac{\left(\gamma_{V}^{2}-1\right) \cdot a_{x 2} \cdot v_{y 2}}{\gamma_{V} \cdot V}}{a_{y 2}=} \frac{\left\{\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{x 2}}{\gamma_{V} \cdot V}+\gamma_{V}\right]^{3}}{\left\{a _ { y 1 } \cdot \left[\frac{\left(1-\gamma_{V}^{2}\right) \cdot v_{x 1}}{\left.\left.\gamma_{V} \cdot \gamma_{V}\right]\right\}-\frac{\left(1-\gamma_{V}^{2}\right) \cdot a_{x 1} \cdot v_{y 1}}{\gamma_{V} \cdot V}}\right.\right.}  \tag{15}\\
a_{z 1}=\frac{\left[\frac{\left(1-\gamma_{V}^{2}\right) \cdot v_{x 1}}{\gamma_{V} \cdot V}+\gamma_{V}\right]^{3}}{\left\{a _ { z 2 } \cdot \left[\frac{\left.\left.\left(\gamma_{V}^{2}-1\right) \cdot v_{x 2}+\gamma_{V}\right]\right\}-\frac{\left(\gamma_{V}^{2}-1\right) \cdot a_{x 2} \cdot v_{z 2}}{\gamma_{V} \cdot V}}{\gamma_{V} \cdot V}\right.\right.}  \tag{16}\\
a_{z 2}=\frac{\left[\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{x 2}}{\gamma_{V} \cdot V}+\gamma_{V}\right]^{3}}{\left\{a _ { z 1 } \cdot \left[\frac{\left.\left.\left(1-\gamma_{V}^{2}\right) \cdot v_{x 1}+\gamma_{V}\right]\right\}-\frac{\left(1-\gamma_{V}^{2}\right) \cdot a_{x 1} \cdot v_{z 1}}{\gamma_{V} \cdot V}}{\left[\frac{\left(1-\gamma_{V}^{2}\right) \cdot v_{x 1}}{\gamma_{V} \cdot V}+\gamma_{V}\right]^{3}}\right.\right.} \tag{17}
\end{gather*}
$$

## 3. The equation of contact for the coefficients of proportionality

Consider three inertial reference systems, as shown in Fig.2, stationary $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and mobile $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ and $\mathrm{O}_{3} x_{3} y_{3} z_{3}$, in which:

- similar the axis of the Cartesian coordinate systems $\mathrm{O}_{1} x_{1} y_{1} z_{1}, \mathrm{O}_{2} x_{2} y_{2} z_{2}$ and $\mathrm{O}_{3} x_{3} y_{3} z_{3}$ are parallel to the three and equally directed;
- system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ moves relative to the system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ with constant speed $V_{2}$ along the axis $\mathrm{O}_{1} x_{1}$;
- system $\mathrm{O}_{3} x_{3} y_{3} z_{3}$ moves relative to the system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ with constant speed $V_{3}$ along the axis $\mathrm{O}_{1} x_{1}$;
- in these three systems as the start timing ( $t_{1}=0, t_{2}=0$ and $\left.t_{3}=0\right)$ is selected,
when their origin $\mathrm{O}_{1}, \mathrm{O}_{2}$ and $\mathrm{O}_{3}$ are the same.


Fig. 2
Based on the formula (8), we can determine the value of the speed $V_{23}$ of the point $\mathrm{O}_{3}$ relative to the point $\mathrm{O}_{2}$ :

$$
\begin{equation*}
V_{23}=\frac{V_{3}-V_{2}}{\frac{\left(1-\gamma_{V 2}^{2}\right) \cdot V_{3}}{\gamma_{V 2}^{2} \cdot V_{2}}+1} \tag{19}
\end{equation*}
$$

and the value of the speed $V_{32}$ of the point $\mathrm{O}_{2}$ relative to the point $\mathrm{O}_{3}$ :

$$
\begin{equation*}
V_{32}=\frac{V_{2}-V_{3}}{\frac{\left(1-\gamma_{V 3}^{2}\right) \cdot V_{2}}{\gamma_{V 3}^{2} \cdot V_{3}}+1} \tag{20}
\end{equation*}
$$

where: $\gamma_{V 2}$ and $\gamma_{V 3}$-coeffic ients of proportionality for inertial reference systems moving relative to a stationary reference system at the speed $V_{2}$ and $V_{3}$, respectively.

Using the principle of relativity, by which point $\mathrm{O}_{3}$ will be removed on the point $\mathrm{O}_{2}$ at the speed equal in magnitude and oppositely directed of the speed, with which the point $\mathrm{O}_{2}$ is removed relative to the point $\mathrm{O}_{3}$, ie:

$$
\begin{equation*}
V_{32}=-V_{23} \tag{21}
\end{equation*}
$$

Substituting equation (21) into formulas (19) and (20), we obtain:

$$
\begin{equation*}
\frac{\left(1-\gamma_{V 2}^{2}\right) \cdot V_{3}}{\gamma_{V 2}^{2} \cdot V_{2}}+1=\frac{\left(1-\gamma_{V 3}^{2}\right) \cdot V_{2}}{\gamma_{V 3}^{2} \cdot V_{3}}+1 \tag{22}
\end{equation*}
$$

From equation (22), it follows that:

$$
\begin{equation*}
\frac{\gamma_{V 2}^{2}-1}{\gamma_{V 2}^{2} \cdot V_{2}^{2}}=\frac{\gamma_{V 3}^{2}-1}{\gamma_{V 3}^{2} \cdot V_{3}^{2}} \tag{23}
\end{equation*}
$$

Since the values of the coefficients of proportionality $\gamma_{V 2}$ and $\gamma_{V 3}$ not depend on each other and depend only on the values of the speeds $V_{2}$ and $V_{3}$, respectively, and the speeds $V_{2}$ and $V_{3}$ were set arbitrarily (and do not depend on each other), then we can say that:

$$
\begin{equation*}
\frac{\gamma_{V 2}^{2}-1}{\gamma_{V 2}^{2} \cdot V_{2}^{2}}=\frac{\gamma_{V 3}^{2}-1}{\gamma_{V 3}^{2} \cdot V_{3}^{2}}=K=\text { Const } \tag{24}
\end{equation*}
$$

ie obtained in general terms that:

$$
\begin{equation*}
\frac{\gamma_{V}^{2}-1}{\gamma_{V}^{2} \cdot V^{2}}=K=\text { Const } \tag{25}
\end{equation*}
$$

where: $K$ - constant, irrespective of the magnitude of the speed $V$ and the coefficient of proportionality $\gamma_{V}$.

As seen from formula (25), depending on the value of the coefficient of proportionality $\gamma_{V}$ the constant $K$ may have the following meanings:

- with $\gamma_{V}=1$ the constant $K$ will be equal to 0 ;
- if the coefficient of proportionality $\gamma_{V}>1$, then the constant $K$ will be positive, ie $K>0$;
- if the coefficient of proportionality $0<\gamma_{V}<1$, then the constant $K$ will have a negative value, ie $K<0$.

From equation (25) can obtain a formula for the coefficient of proportionality $\gamma_{V}$ :

$$
\begin{equation*}
\gamma_{V}^{2}=\frac{1}{1-\left(K \cdot V^{2}\right)} \tag{26}
\end{equation*}
$$

By analogy with the special theory of relativity, we assume that:

- for values of the coefficient of proportionality $\gamma_{V}>1$ the constant $K$ is equal to:

$$
\begin{equation*}
K=\frac{1}{c_{1}^{2}} \tag{27}
\end{equation*}
$$

- for values of the coefficient of proportionality $0<\gamma_{V}<1$ the constant $K$ is equal to:

$$
\begin{equation*}
K=-\frac{1}{c_{2}^{2}} \tag{28}
\end{equation*}
$$

where: $c_{1}$ and $c_{2}$ - real constants.

## 4. Determination of specific speed

Suppose that in a stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ there exists a value $v_{x k r}$ of the speed projection $v_{x 1}$ of point A, shown in Fig. 1, which in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ would be consistent value of the speed projection $v_{x 2}$ of the same point A, equal $v_{x k r}$, ie .:

$$
\begin{equation*}
v_{x 1}=v_{x 2}=v_{x k r} \tag{29}
\end{equation*}
$$

We call $v_{x k r}$ specific speed.
Substituting (29) into formula (7) or (8), we obtain the dependence of the specific speed $v_{x k r}$ on the magnitude of the speed $V$ and the coefficient of proportionality $\gamma_{V}$ :

$$
\begin{equation*}
v_{x k r}^{2}=\frac{\gamma_{V}^{2} \cdot \mathrm{~V}^{2}}{\gamma_{V}^{2}-1} \tag{30}
\end{equation*}
$$

As seen from formula (30):

- for values of the coefficient of proportionality $\gamma_{V}$, which are in the range $\gamma_{V}>1$, equity of the projections $v_{x 1}$ and $v_{x 2}$ of speeds is possible, because when $\gamma_{V}>1$ the specific speed $v_{x k r}$ will have real meaning;
- for values of the coefficient of proportionality $\gamma_{V}$, which are in the range $0<\gamma_{V}<1$, the equality of the projections $v_{x 1}$ and $v_{x 2}$ of speeds is not possible, ie value $v_{x 1}$ can never be equal to the value $v_{x 2}$, because when $0<\gamma_{V}<1$ the specific speed $v_{x k r}$ will have an imaginary value.

From formula (30) can be obtained dependence of the coefficient of proportionality $\gamma_{V}$ on the magnitude of the speed $V$ :

$$
\begin{equation*}
\gamma_{V}^{2}=\frac{1}{1-\frac{V^{2}}{v_{x k r}^{2}}} \tag{31}
\end{equation*}
$$

If we return to formula (26) and compare it with the formula (31), we can
see that:

$$
\begin{equation*}
K=\frac{1}{v_{x k r}^{2}} \tag{32}
\end{equation*}
$$

Based on the fact that the constant $K$ is a constant, it may be noted that $v_{x k r}^{2}$ will be constant, independent of the values of the speed $V$ and the coefficient of proportionality $\gamma_{V}$.

## 5. Basic kinematic equations for the values of the coefficient of proportionality $\gamma_{V}$ in the range $\gamma_{V}>1$ and $0<\gamma_{V}<1$

Using formula (31) taking into account equation (27) for the coefficient of proportionality $\gamma_{V}$, which has the values $\gamma_{V}>1$, which is denoted as $\gamma_{V}>$, we can write:

$$
\begin{equation*}
\gamma_{V}^{2}>=\frac{1}{1-\frac{V^{2}}{c_{1}^{2}}} \tag{33}
\end{equation*}
$$

And from the formula (31) taking into account equation (28) for the coefficient of proportionality $\gamma_{V}$, having the values $0<\gamma_{V}<1$, which is denoted as $\gamma_{V<}$, we can get:

$$
\begin{equation*}
\gamma_{V<}^{2}=\frac{1}{1+\frac{V^{2}}{c_{2}^{2}}} \tag{34}
\end{equation*}
$$

Equation, similar to formula (34), was obtained Y.P. Terletsky [4] and discarded them without a theoretical justification as contrary to experience.

Substituting formulas (33) and (34) in equation (1) - (14), we obtain two systems of equations, which are located opposite each other for comparison, and the sign «<»» indicates that this is the case when $\gamma_{V}>1$, and sign «<»» - for the case when $0<\gamma_{V}<1$ :

$$
\begin{equation*}
\left.x_{1>}=\frac{x_{2>}+\left(V \cdot t_{2>}\right)}{\sqrt{1-\frac{V^{2}}{c_{1}^{2}}}} \quad(35) \right\rvert\, x_{1<}=\frac{x_{2<}+\left(V \cdot t_{2<}\right)}{\sqrt{1+\frac{V^{2}}{c_{2}^{2}}}} \tag{49}
\end{equation*}
$$

$$
\begin{align*}
& x_{2>}=\frac{x_{1>}-\left(V \cdot t_{1>}\right)}{\sqrt{1-\frac{V^{2}}{c_{1}^{2}}}}  \tag{36}\\
& y_{1>}=y_{2>} \\
& z_{1>}=z_{2>}  \tag{52}\\
& t_{1>}=\frac{t_{2>}+\frac{V \cdot x_{2>}}{c_{1}^{2}}}{\sqrt{1-\frac{V^{2}}{c_{1}^{2}}}}  \tag{53}\\
& t_{2>}=\frac{t_{1>}-\frac{V \cdot x_{1>}}{c_{1}^{2}}}{\sqrt{1-\frac{V^{2}}{c_{1}^{2}}}}  \tag{54}\\
& v_{x 1>}=\frac{v_{x 2>}+V}{1+\frac{V \cdot v_{x 2>}}{c_{1}^{2}}}  \tag{55}\\
& v_{x 2>}=\frac{v_{x 1>}-V}{1-\frac{V \cdot v_{x 1>}}{c_{1}^{2}}}  \tag{56}\\
& v_{y 1>}=\frac{v_{y 2>} \cdot \sqrt{1-\frac{V^{2}}{c_{1}^{2}}}}{1+\frac{V \cdot v_{x 2>}}{c_{1}^{2}}}  \tag{57}\\
& v_{y 2>}=\frac{v_{y 1>} \cdot \sqrt{1-\frac{V^{2}}{c_{1}^{2}}}}{1-\frac{V \cdot v_{x 1>}}{c_{1}^{2}}}  \tag{58}\\
& v_{z 1>}=\frac{v_{z 2>} \cdot \sqrt{1-\frac{V^{2}}{c_{1}^{2}}}}{1+\frac{V \cdot v_{x 2>}}{c_{1}^{2}}}  \tag{59}\\
& x_{2<}=\frac{x_{1<}-\left(V \cdot t_{1<}\right)}{\sqrt{1+\frac{V^{2}}{c_{2}^{2}}}}  \tag{50}\\
& y_{1<}=y_{2<}  \tag{51}\\
& z_{1<}=Z_{2<} \\
& t_{1<}=\frac{t_{2<}-\frac{V \cdot x_{2<}}{c_{2}^{2}}}{\sqrt{1+\frac{V^{2}}{c_{2}^{2}}}} \\
& t_{2<}=\frac{t_{1<}+\frac{V \cdot x_{1<}}{c_{2}^{2}}}{\sqrt{1+\frac{V^{2}}{c_{2}^{2}}}} \\
& v_{x 1<}=\frac{v_{x 2<}+V}{1-\frac{V \cdot v_{x 2<}}{c_{2}^{2}}} \\
& v_{x 2<}=\frac{v_{x 1<}-V}{1+\frac{V \cdot v_{x 1<}}{c_{2}^{2}}} \\
& v_{y 1<}=\frac{v_{y 2<} \cdot \sqrt{1+\frac{V^{2}}{c_{2}^{2}}}}{1-\frac{V \cdot v_{x 2<}}{c_{2}^{2}}} \\
& v_{y 2<}=\frac{v_{y 1<} \cdot \sqrt{1+\frac{V^{2}}{c_{2}^{2}}}}{1+\frac{V \cdot v_{x 1<}}{c_{2}^{2}}} \\
& v_{z 1<}=\frac{v_{z 2<} \cdot \sqrt{1+\frac{V^{2}}{c_{2}^{2}}}}{1-\frac{V \cdot v_{x 2<}}{c_{2}^{2}}}
\end{align*}
$$

$$
\begin{align*}
& v_{z 2>}=\frac{v_{z 1>} \cdot \sqrt{1-\frac{V^{2}}{c_{1}^{2}}}}{1-\frac{V \cdot v_{x 1>}}{c_{1}^{2}}}(46) \quad v_{z 2<}=\frac{v_{z 1<} \cdot \sqrt{1+\frac{V^{2}}{c_{2}^{2}}}}{1+\frac{V \cdot v_{x 1<}}{c_{2}^{2}}}  \tag{60}\\
& a_{x 1>}=  \tag{61}\\
& a_{x 2>} \cdot\left(\sqrt{1-\frac{V^{2}}{c_{1}^{2}}}\right)^{3}  \tag{62}\\
& \left(1+\frac{V \cdot v_{x 2>}}{c_{1}^{2}}\right)^{3}
\end{align*}(47) \quad a_{x 1<}=\frac{a_{x 2<} \cdot\left(\sqrt{1+\frac{V^{2}}{c_{2}^{2}}}\right)^{3}}{\left(1-\frac{V \cdot v_{x 2<}}{c_{2}^{2}}\right)^{3}}
$$

For convenience of comparison the above formulas, you can use the following graphs:

- graph, shown in Fig.3, of the length of the segment $\Delta x_{1}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$, which in the mobile inertial reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ corresponds to the segment $\Delta x_{2}$, having fixed ends, on the speed $V$ :


Fig. 3

- graph, shown in Fig.4, of the time interval $\Delta t_{1}$ between two events in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$, which in the mobile inertial reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ occurred in the time interval $\Delta t_{2}$ in the same point on the speed V:


Fig. 4

- graph, shown in Fig.5, of the relationship between the projection $v_{x 2}$ of the speed of the point in the mobile inertial reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ and the projection $v_{x 1}$ of the speed of this point in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ (at constant speed $V$ ):


Fig. 5

- graph, shown in Fig.6, of the relationship between the projection $v_{x 1}$ of the speed of the point in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and the projection $v_{x 2}$ of the speed of this point in the mobile inertial reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ (at constant speed $V$ ):


Fig. 6

- graph, shown in Fig.7, of the relationship between the projection $a_{x 2}$ of the acceleration of the point in the mobile inertial reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ and the projection $v_{x 1}$ of the speed of this point in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ (at constant speed $V$ and constant projection $a_{x 1}$ of the acceleration):


Fig. 7

- graph, shown in Fig.8, of the relationship between the projection $a_{x 1}$ of the acceleration of the point in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and the projection $v_{x 2}$ of the speed of this point in the mobile inertial reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ (at constant speed $V$ and constant projection $a_{x 2}$ of the acceleration):


Fig. 8

Also be noted that all physical processes occur in the four-dimensional space-time, whose geometry - is pseudoeuclidean and determined invariant $J=\left(c_{1}^{2} \cdot t^{2}\right)-x^{2}-y^{2}-z^{2} \quad[12]$ for the case when the coefficient of proportionality $\gamma_{V}>1$, and the invariant $J=\left(c_{2}^{2} \cdot t^{2}\right)+x^{2}+y^{2}+z^{2}$ for the case when the coefficient of proportionality $0<\gamma_{V}<1$.

## 6. Dependences of mass, momentum and kinetic energy of the moving body from its speed

Based on the mandatory implementation of the laws of conservation of momentum and energy of a closed mechanical system, the dependence of mass $M(v)$ of a moving body on the speed $v$ can be obtained using the Lagrangian [1], [8], [9], [12], [14], [23], when considering the perfectly elastic or perfectly plastic collision [15], [16], [17], [18], [19], [20], [21], or just intuitively [13],
[22].
Also, the dependence of mass $M(v)$ of the moving the body on the speed v can be obtained by selecting the function of this dependence in the equations, written for two inertial reference systems, based on the laws of conservation of momentum and energy of a closed mechanical system consisting of two bodies facing perfectly elastic direct central collision, bearing short-term nature, with different positions of the system of bodies in space [28].

Summarizing the results of the findings [8], [13], [22], [28], the dependence of mass $M(v)$ of the moving body, having a rest mass $M_{\mathrm{o}}$, on the speed $v$ is as follows:

$$
\begin{equation*}
M(v)=M_{0} \cdot \gamma_{v} \tag{63}
\end{equation*}
$$

where: $\gamma_{v}$-coefficient of proportionality with the speed $V$, equal to $v$.
Knowing the relationship [1], [28] between the mass of the moving body and its momentum $P(v)$ and the kinetic energy $E_{k i n}(v)$, we can write:

$$
\begin{align*}
& P(v)=M_{0} \cdot \gamma_{v} \cdot v  \tag{64}\\
& E_{k i n}(v)=\frac{\mathrm{M}_{0} \cdot \gamma_{v}^{2} \cdot v^{2}}{\gamma_{v}+1} \tag{65}
\end{align*}
$$

Check the correctness of the choice of formulas (63) - (65) can be in the following example 1.

Assume that there are two inertial reference systems, similar system of reference, shown in Fig.1, stationary $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and mobile $\mathrm{O}_{2} x_{2} y_{2} z_{2}$, which moves with speed $V$ parallel to the axis $\mathrm{O}_{1} x_{1}$ on the system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$.

Suppose that there is a closed mechanical system, consisting of a body 1 and body 2 (as shown in Fig.9), with a masses at rest, equal $M_{\mathrm{o} 1}$ and $M_{\mathrm{o} 2}$, respectively.


Fig. 9

In the mobile reference system $O_{2} x_{2} y_{2} z_{2}$ to a certain point in time $t_{2 c}$ body 1 and body 2 moving parallel to the axis $O_{2} x_{2}$ on one line with constant speeds $v_{21 x n}$ and $v_{22 x n}$ respectively, ie to the time, smaller $t_{2 c}$, body 1 had the momentum $P_{21 x n}$ and the kinetic energy $E_{\text {kin 21xn }}$, and the body 2 had the momentum $P_{22 x n}$ and the kinetic energy $E_{k i n} 22 x n$.

At some point in time $t_{2 c}$ between bodies 1 and 2, there was a direct central perfectly elastic collis ion.

Then, after the collision, at the time, more $t_{2 c}$, bodies 1 and 2 are moving parallel to the axis $O_{2} x_{2}$ on one line with constant speeds $v_{21 x k}$ and $v_{22 x k}$ respectively, ie at the time, greater $t_{2 c}$, body 1 had the momentum $P_{21 x k}$ and the kinetic energy $E_{k i n 21 x k}$, and the body 2 had the momentum $P_{22 x k}$ and the kinetic energy $E_{k i n 22 x k}$.

In the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ collision between bodies 1 and 2 occurred at the time $t_{1 c}$, corresponding to time $t_{2 c}$ in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$.

In the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ to a certain point in time $t_{1 c}$ body 1 and body 2 moving parallel to the axis $\mathrm{O}_{1} x_{1}$ on one line with constant speeds $v_{11 x n}$ and $v_{12 x n}$, respectively, ie to the time, smaller $t_{1 c}$, body 1 had the momentum $P_{11 x n}$ and the kinetic energy $E_{k i n 11 x n}$, and the body 2 had the momentum $P_{12 x n}$ and the kinetic energy $E_{\text {kin } 12 x n}$.

Then, after the collision, at the time, more $t_{1 c}$, bodies 1 and 2 are moving
parallel to the axis $\mathrm{O}_{1} x_{1}$ on one line with constant speeds $v_{11 x k}$ and $v_{12 x k}$ respectively, ie at the time, greater $t_{1 c}$, body 1 had the momentum $P_{11 x k}$ and the kinetic energy $E_{k i n 11 x k}$, and the body 2 had the momentum $P_{12 x k}$ and the kinetic energy $E_{k i n 12 x k}$.

Given that:

- has the symmetry of space and time,
- body 1 and body 2 form the closed mechanical system,
- between bodies 1 and 2 there was a direct central collision,
- collision between bodies 1 and 2 was of an elastic nature,
we can write the conservation laws of momentum and mechanical energy for the closed mechanical system, consisting of bodies 1 and 2 , considering the moments before and after the collision:
in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ :

$$
\begin{gather*}
P_{11 x n}+P_{12 x n}=P_{11 x k}+P_{12 x k}  \tag{66}\\
E_{k i n ~ 11 x n}+E_{k i n 12 x n}=E_{k i n 11 x k}+E_{k i n ~ 12 x k} \tag{67}
\end{gather*}
$$

in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ :

$$
\begin{array}{r}
P_{21 x n}+P_{22 x n}=P_{21 x k}+P_{22 x k} \\
E_{\text {kin } 21 x n}+E_{k i n 22 x n}=E_{k i n 21 x k}+E_{k i n 22 x k} \tag{69}
\end{array}
$$

Should also be noted that the speeds $v_{11 x n}$ and $v_{21 x n}, v_{12 x n}$ and $v_{22 x n}, v_{11 x k}$ and $v_{21 x k}, v_{12 x k}$ and $v_{22 x k}$ are interconnected through a conversion speeds (7):

$$
\begin{equation*}
v_{x 11 x n}=\frac{v_{x 21 x n}+V}{\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{x 2 x n}}{\gamma_{V}^{2} \cdot V}+1} \tag{70}
\end{equation*}
$$

etc.
Now, by setting the initial data, we can hold the estimated test of choice of dependences (63) - (65) mass, momentum and kinetic energy of the moving body for the case, when the coefficients of proportionality $\gamma_{V}$ and $\gamma_{v}$ are in the ranges $\gamma_{V}>1$ and $\gamma_{v}>1$, respectively.

Suppose that: $M_{\mathrm{o} 1}=1, M_{\mathrm{o} 2}=0,5, V / c_{1}=0,5, v_{21 \times n} / c_{1}=0,9$, $v_{22 x n} / c_{1}=0,6$.

Then, for this example 1 , with coefficients of proportionality $\gamma_{V}$ and $\gamma_{v}$, whose values may be in the ranges $\gamma_{V}>1$ and $\gamma_{v}>1$, the numerical calculations yield the following results in Tab. 1 for the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ and Tab. 2 for the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$.

Tab. 1
Ranges $\gamma_{V}>1$ and $\gamma_{v}>1$. The mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$.

| Object | Period time | Value | Digit of value |
| :---: | :---: | :---: | :---: |
| Body 1 | Before collision | speed $v_{21 \times n} / c_{1}$ | 0,9 |
|  |  | mass $M_{21 n}$ | 2,294157338706 |
|  |  | momentum $P_{21 \times n} / c_{1}$ | 2,064741604835 |
|  |  | kinetic energy $E_{k i n 21 x n} / c_{1}^{2}$ | 1,294157338706 |
|  | After collision | speed $v_{11 \text { d }} / c_{1}$ | 0,7360143377 |
|  |  | mass $M_{21 k}$ | 1,477179174242 |
|  |  | momentum $P_{21 \times k} / c_{1}$ | 1,087225051595 |
|  |  | kinetic energy $E_{k i n 21 x k} / c_{1}^{2}$ | 0,477179174242 |
| Body 2 | Before collision | speed $v_{22 x n} / c_{1}$ | 0,6 |
|  |  | mass $M_{22 n}$ | 0,625 |
|  |  | momentum $P_{22 \times n} / c_{1}$ | 0,375 |
|  |  | kinetic energy $E_{\text {kin } 22 x n} / c_{1}{ }^{2}$ | 0,125 |
|  | After collision | speed $v_{22 k k} / c_{1}$ | 0,937959108239 |
|  |  | mass $M_{22 k}$ | 1,441978164463 |
|  |  | momentum $P_{22 \times k} / c_{1}$ | 1,35251655324 |
|  |  | kinetic energy $E_{k i n 22 x k} / c_{1}^{2}$ | 0,941978164463 |
| The system of bodies 1 and 2 | Before collision | mass ( $M_{21 n}+M_{22 n}$ ) | 2,919157338706 |
|  |  | $\begin{gathered} \text { momentum } \\ \left(P_{21 \times n}+P_{22 \times n}\right) / c_{1} \end{gathered}$ | 2,439741604835 |
|  |  | $\begin{gathered} \text { kinetic energy } \\ \left(E_{\text {kin } 21 \times n}+E_{\text {kin } 22 x n}\right) / c_{1}^{2} \\ \hline \end{gathered}$ | 1,419157338706 |
|  | After collision | mass ( $M_{21 k}+M_{22 k}$ ) | 2,919157338706 |
|  |  | $\begin{gathered} \text { momentum } \\ \left(P_{21 \times k}+P_{22 x k}\right) / c_{1} \end{gathered}$ | 2,439741604835 |
|  |  | $\left.\begin{array}{c} \text { kinetic energy } \\ \left(E_{k i n}^{21 x k}+E_{k i n}^{22 x k}\right) \end{array}\right) c_{1}^{2}$ | 1,419157338706 |

Tab. 2
Ranges $\gamma_{V}>1$ and $\gamma_{v}>1$. The stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$.

| Object | Period time | Value | Digit of value |
| :---: | :---: | :---: | :---: |
| Body 1 | Before collision | speed $v_{11 x n} / c_{1}$ | 0,965517241379 |
|  |  | mass $M_{11 n}$ | 3,841143835489 |
|  |  | momentum $P_{11 \times n} / c_{1}$ | 3,708690599782 |
|  |  | kinetic energy <br> $E_{\text {kin } 11 x n} / c_{1}{ }^{2}$ | 2,841143835489 |
|  | After collision | speed $v_{11 x k} / c_{1}$ | 0,903514517939 |
|  |  | mass $M_{11 k}$ | 2,333409263988 |
|  |  | momentum $P_{11 x k} / c_{1}$ | 2,108269146306 |
|  |  | kinetic energy <br> $E_{k i n 11 x k} / c_{1}^{2}$ | 1,333409263988 |
| Body 2 | Before collision | speed $v_{12 \times n} / c_{1}$ | 0,846153846154 |
|  |  | mass $M_{12 n}$ | 0,938194187433 |
|  |  | momentum $P_{12 x n} / c_{1}$ | 0,793856620136 |
|  |  | kinetic energy <br> $E_{\text {kin } 12 x n} / c_{1}^{2}$ | 0,438194187433 |
|  | After collis ion | speed $v_{12 x k} / c_{1}$ | 0,978882996844 |
|  |  | mass $M_{12 k}$ | 2,445928758933 |
|  |  | momentum $P_{12 x k} / c_{1}$ | 2,394278073612 |
|  |  | kinetic energy $E_{k i n 12 x k} / c_{1}^{2}$ | 1,945928758933 |
| The system of bodies 1 and 2 | Before collision | mass $\left(M_{11 n}+M_{12 n}\right)$ | 4,779338022922 |
|  |  | $\begin{gathered} \text { momentum } \\ \left(P_{11 \times n}+P_{12 x n}\right) / c_{1} \end{gathered}$ | 4,502547219918 |
|  |  | $\begin{gathered} \text { kinetic energy } \\ \left(E_{\text {kin } 11 x n}+E_{\text {kin } 12 x n}\right) / c_{1}^{2} \end{gathered}$ | 3,279338022922 |
|  | After collision | mass ( $M_{11 k}+M_{12 k}$ ) | 4,779338022922 |
|  |  | $\begin{gathered} \text { momentum } \\ \left(P_{11 x k}+P_{12 x k}\right) / c_{1} \end{gathered}$ | 4,502547219918 |
|  |  | $\begin{gathered} \text { kinetic energy } \\ \left(E_{\text {kin } 11 x k}+E_{\text {kin } 12 x k}\right) / c_{1}^{2} \end{gathered}$ | 3,279338022922 |

As a result of the calculation can draw the following conclusion: in example 1 in the mobile $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ and stationary $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ reference systems before and after the collision the mass, momentum and kinetic energy of the mechanical system of bodies 1 and 2 remain unchanged for the case, when the coefficients of proportionality $\gamma_{V}$ and $\gamma_{v}$ lie in the ranges $\gamma_{V}>1$ and $\gamma_{v}>1$, when using the relationships (63) - (65).

Also, by setting the initial data, we can hold the estimated test of choice of dependences (63) - (65) mass, momentum and kinetic energy of the moving body for the case, when the coefficients of proportionality $\gamma_{V}$ and $\gamma_{v}$ are in the ranges $0<\gamma_{V}<1$ and $0<\gamma_{v}<1$, respectively.

Suppose that: $M_{\mathrm{o} 1}=1, M_{\mathrm{o} 2}=0,5, V / c_{2}=0,5, v_{21 \times n} / c_{2}=0,9$, $v_{22 x n} / c_{2}=0,6$.

For this example 1 , with coefficients of proportionality $\gamma_{V}$ and $\gamma_{v}$, whose values may be in the ranges $0<\gamma_{V}<1$ and $0<\gamma_{v}<1$, the numerical calculations yield the following results in Tab. 3 for the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ and Tab. 4 for the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$.

Tab. 3
Ranges $0<\gamma_{V}<1$ and $0<\gamma_{v}<1$. The mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$.

| Object | Period time | Value | Digit of value |
| :---: | :---: | :---: | :---: |
| Body 1 | Before collision | speed $v_{21 \times n} / c_{2}$ | 0,9 |
|  |  | mass $M_{21 n}$ | 0,743294146247 |
|  |  | momentum $P_{21 \times n} / c_{2}$ | 0,668964731622 |
|  |  | kinetic energy $E_{\text {kin } 21 x n} / c_{2}^{2}$ | 0,256705853753 |
|  | After collision | speed $v_{11 x k} / c_{2}$ | 0,691099932748 |
|  |  | mass $M_{21 k}$ | 0,822656908881 |
|  |  | momentum $P_{21 x k} / c_{2}$ | 0,568538134403 |
|  |  | kinetic energy $E_{k i n 21 x k} / c_{2}^{2}$ | 0,177343091119 |
| Body 2 | Before collision | speed $v_{22 \times n} / c_{2}$ | 0,6 |
|  |  | mass $M_{22 n}$ | 0,428746462856 |
|  |  | momentum $P_{22 x n} / c_{2}$ | 0,257247877714 |
|  |  | kinetic energy $E_{\text {kin } 22 x n} / c_{2}{ }^{2}$ | 0,071253537144 |
|  | After collision | speed $v_{22 x k} / c_{2}$ | 1,023729712365 |
|  |  | mass $M_{22 k}$ | 0,349383700222 |
|  |  | momentum $P_{22 x k} / c_{2}$ | 0,357674474934 |
|  |  | kinetic energy $E_{k i n 22 x k} / c_{2}^{2}$ | 0,150616299778 |
| The system of bodies 1 and 2 | Before collision | mass ( $M_{21 n}+M_{22 n}$ ) | 1,172040609103 |
|  |  | $\begin{gathered} \text { momentum } \\ \left(P_{21 x n}+P_{22 x n}\right) / c_{2} \end{gathered}$ | 0,926212609336 |
|  |  | $\begin{gathered} \text { kinetic energy } \\ \left(E_{\text {kin } 21 x n}+E_{\text {kin } 22 x n}\right) / c_{2}^{2} \end{gathered}$ | 0,327959390897 |
|  | After collision | mass $\left(M_{21 k}+M_{22 k}\right)$ | 1,172040609103 |
|  |  | momentum $\left(P_{21 x k}+P_{22 x k}\right) / c_{2}$ | 0,926212609336 |
|  |  | $\begin{gathered} \text { kinetic energy } \\ \left(E_{\text {kin } 21 x k}+E_{\text {kin } 22 x k}\right) / c_{2}^{2} \end{gathered}$ | 0,327959390897 |

Tab. 4
Ranges $0<\gamma_{V}<1$ and $0<\gamma_{v}<1$. The stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$.

| Object | Period time | Value | Digit of value |
| :---: | :---: | :---: | :---: |
| Body 1 | Before collision | speed $v_{11 x n} / c_{2}$ | 2,545454545455 |
|  |  | mass $M_{11 n}$ | 0,365652372423 |
|  |  | momentum $P_{11 \times n} / c_{2}$ | 0,93075149344 |
|  |  | kinetic energy $E_{\text {kin } 11 x n} / c_{2}{ }^{2}$ | 0,634347627577 |
|  | After collision | speed $v_{11 x k} / c_{2}$ | 1,820001331727 |
|  |  | mass $M_{11 k}$ | 0,481548724902 |
|  |  | momentum $P_{11 x k} / c_{2}$ | 0,876419320614 |
|  |  | kinetic energy $E_{\text {kin } 11 x k} / c_{2}{ }^{2}$ | 0,518451275098 |
| Body 2 | Before collision | speed $v_{12 \times n} / c_{2}$ | 1,571428571429 |
|  |  | mass $M_{12 n}$ | 0,268437746097 |
|  |  | momentum $P_{12 \times n} / c_{2}$ | 0,421830743866 |
|  |  | kinetic energy <br> $E_{\text {kin } 12 x n} / c_{2}{ }^{2}$ | 0,231562253903 |
|  | After collision | speed $v_{12 x k} / c_{2}$ | 3,121532492927 |
|  |  | mass $M_{12 k}$ | 0,152541393617 |
|  |  | momentum $P_{12 x k} / c_{2}$ | 0,476162916693 |
|  |  | kinetic energy $E_{k i n 12 x k} / c_{2}^{2}$ | 0,347458606383 |
| The system of bodies 1 and 2 | Before collision | mass ( $M_{11 n}+M_{12 n}$ ) | 0,63409011852 |
|  |  | $\begin{gathered} \text { momentum } \\ \left(P_{11 x n}+P_{12 x n}\right) / c_{2} \end{gathered}$ | 1,352582237306 |
|  |  | $\begin{gathered} \text { kinetic energy } \\ \left(E_{\text {kin } 11 x n}+E_{\text {kin } 12 x n}\right) / c_{2}^{2} \end{gathered}$ | 0,86590988148 |
|  | After collision | mass ( $M_{11 k}+M_{12 k}$ ) | 0,63409011852 |
|  |  | $\begin{gathered} \text { momentum } \\ \left(P_{11 x k}+P_{12 x k}\right) / c_{2} \end{gathered}$ | 1,352582237306 |
|  |  | $\begin{gathered} \text { kinetic energy } \\ \left(E_{\text {kin } 11 x k}+E_{\text {kin } 12 x k}\right) / c_{2}^{2} \end{gathered}$ | 0,86590988148 |

Here also as a result of the calculation can draw the following conclusion:
in example 1 in the mobile $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ and stationary $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ reference systems before and after the collision the mass, momentum and kinetic energy of the mechanical system of bodies 1 and 2 remain unchanged for the case, when the coefficients of proportionality $\gamma_{V}$ and $\gamma_{v}$ lie in the ranges $0<\gamma_{V}<1$ and $0<\gamma_{V}<1$, when using the relationships (63)-(65).

Substituting formulas (33) and (34) in equations (63) - (65), we can obtain the dependences of mass, momentum and kinetic energy of the moving body on its speed for cases, where the coefficient of proportionality $\gamma_{v}>1$ and $0<\gamma_{v}<1$, which are located opposite each other for comparison, and the sign «>» indicates that this is the case when $\gamma_{v}>1$, and sign «<» - for the case when $0<\gamma_{v}<1$ :

$$
\begin{align*}
& M(v)_{>}=\frac{M_{o}}{\sqrt{1-\frac{v^{2}}{c_{1}^{2}}}}  \tag{71}\\
& P(v)_{>}=\frac{M_{0} \cdot v}{\sqrt{1-\frac{v^{2}}{c_{1}^{2}}}}  \tag{72}\\
& =M_{0} \cdot c_{1}^{2} \cdot\left(\frac{1}{\sqrt{1-\frac{v^{2}}{c_{1}^{2}}}}-1\right), ~ \$ ~(v)_{>}= \\
& P(v)_{<}=\frac{M_{0} \cdot v}{\sqrt{1+\frac{v^{2}}{c_{2}^{2}}}} \tag{75}
\end{align*}
$$

Here one can note that a formula similar to formula (74), shows Y.P. Terletsky [4].

The main values of relationships for mass $M(v)_{>}(71)$, momentum $P(v)_{>}$ (72) and kinetic energy $E_{k i n}(v)_{>}$(73) with the coefficient of proportionality
$\gamma_{v}>1$ and dependences of mass $M(v)_{<}(74)$, momentum $P(v)_{<}(75)$ and kinetic energy $E_{\text {kin }}(v)_{<}$(76) with a coefficient of proportionality $0<\gamma_{v}<1$ are given in Tab. 5 and Tab.6, respectively:

Tab. 5
For the coefficient of proportionality $\gamma_{v}>1$

| Speed $v$ | Mass $M(v)_{>}$ | Momentum $P(v)_{>}$ | Kinetic energy <br> $E_{\text {kin }}(v)_{>}$ |
| :---: | :---: | :---: | :---: |
| $v \ll c_{1}$ | $M_{0}$ | $M_{0} \cdot v$ | $\frac{M_{0} \cdot v^{2}}{2}$ |
| $v<c_{1}$ | has real meaning | has real meaning | has real meaning |
| $v=c_{1}$ | $\infty$ | $\infty$ | $\infty$ |
| $v>c_{1}$ | has no real meaning | has no real <br> meaning | has no real <br> meaning |

Tab. 6
For the coefficient of proportionality $0<\gamma_{v}<1$

| Speed $v$ | Mass $M(v)_{<}$ | Momentum $P(v)_{<}$ | Kinetic energy <br> $E_{k i n}(v)_{<}$ |
| :---: | :---: | :---: | :---: |
| $v \ll c_{2}$ | $M_{\mathrm{o}}$ | $M_{\mathrm{o}} \cdot v$ | $\frac{M_{\mathrm{o}} \cdot v^{2}}{2}$ |
| $v<c_{2}$ | has real meaning | has real meaning | has real meaning |
| $v=c_{2}$ | $\frac{M_{\mathrm{o}}}{\sqrt{2}}$ | $\frac{M_{\mathrm{o}} \cdot c_{2}}{\sqrt{2}}$ | $M_{\mathrm{o}} \cdot c_{2}^{2} \cdot\left(1-\frac{1}{\sqrt{2}}\right)$ |
| $v>c_{2}$ | has real meaning | has real meaning | has real meaning |
| $v=\infty$ | tends to zero | $M_{\mathrm{o}} \cdot c_{2}$ | $M_{\mathrm{o}} \cdot c_{2}^{2}$ |

As seen from tab. 5 and tab. 6 , both the range of possible values of the coefficient of proportionality $\gamma_{v}>1$ and $0<\gamma_{v}<1$ are equivalent (both satisfy the boundary condition).

For comparison of formulas (71) - (73) and (74) - (76) give the following graphs:

- graphs of dependence of the mass $M(v)$ of a moving body on the speed $v$, shown in Fig.10:


Fig. 10

- graphs of dependence of the momentum $P(v)$ of a moving body on the speed $v$, shown in Fig.11:


Fig. 11

- graphs of dependence of the kinetic energy $E_{\text {kin }}(v)$ of a moving body on the speed $v$, shown in Fig. 12:


Fig. 12

## 7. Defining values constants $\boldsymbol{c}_{1}$ and $\boldsymbol{c}_{\mathbf{2}}$

The law of conservation of momentum of a closed mechanical system of bodies, connected with the symmetry properties of space - the homogeneity of space [2], states, that the momentum of a closed mechanical system of bodies (which is not acted upon by external forces) is a constant value, ie in any inertial reference system for any point in time the value of the momentum of a closed mechanical system of bodies is a constant value (because there is no external influence).

In the following the above example 2 in the inertial reference system with the help of the special theory of relativity will determine the momentums of the bodies, constituting a closed mechanical system and have been under constant interaction, for two moments of time, then, applying the law of conservation of momentum of a closed mechanical system, will be determined by the values of the constants $c_{1}$ and $c_{2}$.

Assume that there are two inertial reference systems, similar to those of
reference systems, shown in Fig.1, stationary $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and mobile $\mathrm{O}_{2} x_{2} y_{2} z_{2}$, which moves with speed $V$ parallel to the axis $\mathrm{O}_{1} x_{1}$ relative to the system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$.

Suppose that there is a closed mechanical system of bodies, shown in Fig. 13 and consisting of point bodies 1 and 2, with equal mass $M_{\mathrm{o}}$ at rest, and a string 3 .


Fig. 13

Bodies 1 and 2 are connected by a string 3, the mass of which can be neglected because of its smallness.

Bodies 1 and 2 rotate with angular speed $\omega$ around a common center of mass - the point O .

Distance from the point body 1 (body 2 ) to point O is equal to $R$.
Let's put a closed mechanical system of bodies 1 and 2 with a string 3 in the moving reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ so, that the point O would be stationary in this reference system, and coincided with the origin $\mathrm{O}_{2}$, and the rotation of bodies 1 and 2 around it would occur in a clockwise direction in the plane of $\mathrm{O}_{2} x_{2} y_{2}$, as shown in Fig. 14.


Fig. 14

Also assume, that at the start of timing $\left(t_{2}=0\right)$ in the reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ bodies 1 and 2 were on the axis $\mathrm{O}_{2} x_{2}$, with the body 1 had a positive coordinate, and the body 2 - negative.

In the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at any time $t_{2}$ bodies 1 and 2 will have the speeds $v_{21}$ and $v_{22}$, equal $v_{R}$ :

$$
\begin{equation*}
v_{21}=v_{22}=v_{R}=\omega \cdot R \tag{77}
\end{equation*}
$$

In this case, the projections $v_{21 x}$ and $v_{21 y}$ of speed of the body 1 and the projections $v_{22 x}$ and $v_{22 y}$ of speed of the body 2 on the axis $\mathrm{O}_{2} x_{2}$ and $\mathrm{O}_{2} y_{2}$, respectively, for times $t_{21}$ and $t_{22}$ will be equal to:

$$
\begin{gather*}
v_{21 x}=-\left[v_{R} \cdot \sin \left(\omega \cdot t_{21}\right)\right]  \tag{78}\\
v_{21 y}=-\left[v_{R} \cdot \cos \left(\omega \cdot t_{21}\right)\right]  \tag{79}\\
v_{22 x}=v_{R} \cdot \sin \left(\omega \cdot t_{22}\right)  \tag{80}\\
v_{22 y}=v_{R} \cdot \cos \left(\omega \cdot t_{22}\right) \tag{81}
\end{gather*}
$$

The relationship between the coordinates $x_{21}$ and $y_{21}$ of the body 1 depending on time $t_{21}$ and the relationship between the coordinates $x_{22}$ and $y_{22}$ of the body 2 depending on the time $t_{22}$ in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ can be written as:

$$
\begin{gather*}
x_{21}=R \cdot \cos \left(\omega \cdot t_{21}\right)  \tag{82}\\
y_{21}=-\left[R \cdot \sin \left(\omega \cdot t_{21}\right)\right] \tag{83}
\end{gather*}
$$

$$
\begin{gather*}
x_{22}=-\left[R \cdot \cos \left(\omega \cdot t_{22}\right)\right]  \tag{84}\\
y_{22}=R \cdot \sin \left(\omega \cdot t_{22}\right) \tag{85}
\end{gather*}
$$

Based on the equations (1) and (3), we can write the relationships between:

- coordinates $x_{11}$ and $y_{11}$ of the body 1 at time $t_{11}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and coordinates $x_{21}$ and $y_{21}$ of the body 1 in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{21}$, which corresponds to the time $t_{11}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ :

$$
\begin{gather*}
x_{11}=\gamma_{V} \cdot\left[x_{21}+\left(V \cdot t_{21}\right)\right]  \tag{86}\\
y_{11}=y_{21} \tag{87}
\end{gather*}
$$

- coordinates $x_{12}$ and $y_{12}$ of the body 2 at time $t_{12}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and coordinates $x_{22}$ and $y_{22}$ of the body 2 in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{22}$, which corresponds to the time $t_{12}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ :

$$
\begin{gather*}
x_{12}=\gamma_{V} \cdot\left[x_{22}+\left(V \cdot t_{22}\right)\right]  \tag{88}\\
y_{12}=y_{22} \tag{89}
\end{gather*}
$$

Using formula (5) relationship between the values of the times $t_{11}$ and $t_{21}$, $t_{12}$ and $t_{22}$ will look like this:

$$
\begin{align*}
t_{11} & =\frac{\left(\gamma_{V}^{2}-1\right) \cdot x_{21}}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{21}\right)  \tag{90}\\
t_{12} & =\frac{\left(\gamma_{V}^{2}-1\right) \cdot x_{22}}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{22}\right) \tag{91}
\end{align*}
$$

In the considered example 2 , we are interested in the position of bodies 1 and 2 in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ at the same time, ie where:

$$
\begin{equation*}
t_{11}=t_{12} \tag{92}
\end{equation*}
$$

Then equation (92) taking into account formulas (82), (84), (86), (88), (90) and (91) becomes:

$$
\begin{align*}
& \frac{\left(\gamma_{V}^{2}-1\right) \cdot R \cdot \cos \left(\omega \cdot t_{21}\right)}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{21}\right)= \\
= & \frac{\left(1-\gamma_{V}^{2}\right) \cdot R \cdot \cos \left(\omega \cdot t_{22}\right)}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{22}\right) \tag{93}
\end{align*}
$$

In the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ when performing the condition (92) it is interesting position of the bodies 1 and 2 at the time $t_{2 p}$, when:

$$
\begin{equation*}
t_{21}=t_{22}=t_{2 p} \tag{94}
\end{equation*}
$$

Substituting condition (94) in equation (93) for the case when ( $\omega \cdot t_{2 p}$ ) $<\pi$, we obtain:

$$
\begin{equation*}
\omega \cdot t_{2 p}=\frac{\pi}{2} \tag{95}
\end{equation*}
$$

Ie to meet the conditions (92) and (94) during the time $t_{2 p}$ the bodies 1 and 2 should be on a line parallel to the axis $\mathrm{O}_{2} y_{2}$.

Also in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ when performing the condition (92) it is interesting position of body 2 when finding the body 1 on the axis $\mathrm{O}_{2} x_{2}$ at time $t_{21}$, equal to $t_{21 h}$, where:

$$
\begin{equation*}
t_{21 h}=0 \tag{96}
\end{equation*}
$$

The value of time $t_{22}$, when performing the conditions (92) and (96), denote $t_{22 h}$, for which the equation (93) becomes:

$$
\begin{equation*}
t_{22 h}=\left(1-\frac{1}{\gamma_{V}^{2}}\right) \cdot\left[1+\cos \left(\omega \cdot t_{22 h}\right)\right] \cdot \frac{R}{V} \tag{97}
\end{equation*}
$$

or:

$$
\begin{equation*}
\omega \cdot t_{22 h}=\left(1-\frac{1}{\gamma_{V}^{2}}\right) \cdot\left[1+\cos \left(\omega \cdot t_{22 h}\right)\right] \cdot \frac{v_{R}}{V} \tag{98}
\end{equation*}
$$

As seen from equation (98), depending on the value of the coefficient of proportionality $\gamma_{V}$ the value of time $t_{22 h}$ can be:

- $\quad t_{22 h}>0$ when $\gamma_{V}>1$;
- $\quad t_{22 h}<0$ when $0<\gamma_{V}<1$;
- $\quad t_{22 h}=0$ when $\gamma_{V}=1$.

Now we can begin to use the law of conservation of momentum for the preparation of equations.

Consider two points in time in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$.
As a first point in time we choose $t_{1 p}$.
Under the terms of (92), (94) and (95) in the moving mobile reference
system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{2 p}$ the bodies 1 and 2 are on a line parallel to the axis $\mathrm{O}_{2} y_{2}$ and in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ the bodies 1 and 2 will be on a line parallel to the axis $\mathrm{O}_{1} y_{1}$ at time $t_{11}\left(t_{12}\right)$, equal $t_{1 p}$ and which corresponds to the time $t_{2 p}$ in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$.

As shown in Fig.15, according to equations (95), (78) - (81) in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{2 p}$ the bodies 1 and 2, respectively, have the following values of the projections $v_{21 x p}, v_{21 y p}$ and $v_{22 x p}, v_{22 y p}$ of speeds of his movement on the axis $\mathrm{O}_{2} x_{2}$ and $\mathrm{O}_{2} y_{2}$ :

$$
\begin{align*}
& v_{21 x p}=-v_{R}  \tag{99}\\
& v_{21 y p}=0  \tag{100}\\
& v_{22 x p}=v_{R}  \tag{101}\\
& v_{22 y p}=0 \tag{102}
\end{align*}
$$



Fig. 15

Then, on the basis of formulas (7), (9) and equalities (99) - (102), in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ at time $t_{1 p}$ the body 1 and the body 2 , respectively, will have the following values of the projections $v_{11 x p}, v_{11 y p}$ and $v_{12 x p}, v_{12 y p}$ of speeds of his movement on the axis $\mathrm{O}_{1} x_{1}$ and $\mathrm{O}_{1} y_{1}$ :

$$
\begin{gather*}
v_{11 x p}=\frac{V-v_{R}}{1-\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{R}}{\gamma_{V}^{2} \cdot V}}  \tag{103}\\
v_{11 y p}=0  \tag{104}\\
v_{12 x p}=\frac{V+v_{R}}{\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{R}}{\gamma_{V}^{2} \cdot V}+1}  \tag{105}\\
v_{12 y p}=0 \tag{106}
\end{gather*}
$$

Hence, using formula (64), may be noted that in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ at time $t_{1 p}$ the body 1 and the body 2 , respectively, will have the following values of the projections $P_{11 x p}, P_{11 y p}$ and $P_{12 x p}, P_{12 y p}$ of momentums on the axis $\mathrm{O}_{1} x_{1}$ and $\mathrm{O}_{1} y_{1}$ :

$$
\begin{gather*}
P_{11 x p}=M_{\mathrm{o}} \cdot \gamma_{v 11 x \mathrm{p}} \cdot v_{11 x \mathrm{p}}  \tag{107}\\
P_{12 x p}=M_{\mathrm{o}} \cdot \gamma_{v 12 x \mathrm{p}} \cdot v_{12 x \mathrm{p}}  \tag{108}\\
P_{11 y p}=0  \tag{109}\\
P_{12 y p}=0 \tag{110}
\end{gather*}
$$

where: $\gamma_{v 11 x p}$ and $\gamma_{v 12 x p}$ - the coefficients of proportionality of the speed $V$, equal to $v_{11 \times p}$ and $v_{12 x p}$ respectively.

As a second point in time we choose $t_{1 h}$.
Under the terms of (92) and (96) in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{21 h}=0$ the body 1 will be located on the axis $\mathrm{O}_{2} x_{2}$, and in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ the body 1 will be located on the axis $\mathrm{O}_{1} x_{1}$ at time $t_{11}$ $\left(t_{12}\right)$, equal $t_{1 h}$ and which corresponds to the time $t_{21 h}=0$ in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$.

Moreover in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ according to equation (98), when the value of the coefficient of proportionality $\gamma_{V} \neq 1$, the body 2 can not be on the axis $\mathrm{O}_{2} x_{2}$ at time $t_{22 h}$, which corresponds to the time $t_{1 h}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$.

Ie the body 1 is located on the axis $\mathrm{O}_{1} x_{1}$ in a the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ at time $t_{1 h}$, which corresponds to the time $t_{21 h}=0$ in the mobile
reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$, and at the time $t_{1 h}$ the body 2 can not lie on the axis $\mathrm{O}_{2} x_{2}$ (with a coeffic ient of proportionality $\gamma_{V} \neq 1$ ).

As shown in Fig.16, in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ the body 1 at the time $t_{21 h}=0$ and the body 2 at the time $t_{22 h}$ respectively have projections $v_{21 x h}, v_{21 y h}$ and $v_{22 x h}, v_{22 y h}$ of speeds of his movement on the axis $\mathrm{O}_{2} x_{2}$ and $\mathrm{O}_{2} y_{2}$, so that:

$$
\begin{gather*}
v_{21 x h}=0  \tag{111}\\
v_{21 y h}=-v_{R} \tag{112}
\end{gather*}
$$



Fig. 16

Then, on the basis of formulas (7), (9) and equations (111), (112), in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ at time $t_{1 h}$ the body 1 and the body 2 , respectively, will have the values of the projections $v_{11 x h}, v_{11 y h}$ and $v_{12 x h}$, $v_{12 y h}$ of speeds of his movement on the axis $\mathrm{O}_{1} x_{1}$ and $\mathrm{O}_{1} y_{1}$ :

$$
\begin{gather*}
v_{11 x h}=V  \tag{113}\\
v_{11 y h}=-\frac{v_{R}}{\gamma_{V}}  \tag{114}\\
v_{12 x h}=\frac{V+v_{22 x h}}{\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{22 x h}}{\gamma_{V}^{2} \cdot V}+1} \tag{115}
\end{gather*}
$$

$$
\begin{equation*}
v_{12 y h}=\frac{v_{22 y h}}{\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{22 x h}}{\gamma_{V} \cdot V}+\gamma_{V}} \tag{116}
\end{equation*}
$$

Given equation (98), we note that:

- with the coefficient of proportionality $\gamma_{V}>1$ the time $t_{22 h}>0$, so the projection $v_{22 y h}$ of the speed will be the direction of the axis $\mathrm{O}_{2} y_{2}$;
- with the coefficient of proportionality $0<\gamma_{V}<1$ the time $t_{22 h}<0$, so the projection $v_{22 y h}$ of the speed will be the direction opposite to the axis $\mathrm{O}_{2} y_{2}$.

From equations (80) and (81) it follows that:

$$
\begin{equation*}
v_{22 x h}^{2}+v_{22 y h}^{2}=v_{R}^{2} \tag{117}
\end{equation*}
$$

Using formula (64), may be noted that in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ at time $t_{1 h}$ the body 1 and the body 2 , respectively, will have the following values of the projections $P_{11 x h}, P_{11 y h}$ and $P_{12 x h}, P_{12 y h}$ of momentums on the axis $\mathrm{O}_{1} x_{1}$ and $\mathrm{O}_{1} y_{1}$ :

$$
\begin{align*}
& P_{11 x h}=M_{\mathrm{o}} \cdot \gamma_{v 11 h} \cdot v_{11 \mathrm{x} h}  \tag{118}\\
& P_{12 x h}=M_{\mathrm{o}} \cdot \gamma_{v 12 h} \cdot v_{12 \mathrm{x} h}  \tag{119}\\
& P_{11 y h}=M_{\mathrm{o}} \cdot \gamma_{v 11 h} \cdot v_{11 \mathrm{y} h}  \tag{120}\\
& P_{12 y h}=M_{\mathrm{o}} \cdot \gamma_{v 12 h} \cdot v_{12 \mathrm{y} h} \tag{121}
\end{align*}
$$

where: $\gamma_{v 11 h}$ and $\gamma_{v 12 h}$ - the coefficients of proportionality in the speed $V$, equal to $v_{11 h}$ and $v_{12 h}$ respectively, so that:

$$
\begin{align*}
& v_{11 h}^{2}=v_{11 x h}^{2}+v_{11 y h}^{2}  \tag{122}\\
& v_{12 h}^{2}=v_{12 x h}^{2}+v_{12 y h}^{2} \tag{123}
\end{align*}
$$

Due to the fact, that the mechanical system of the bodies 1 and 2 (and string 3 ) is closed, the law of conservation of momentum can write the following equations for the moments of times $t_{1 p}$ and $t_{1 h}$ :

$$
\begin{aligned}
P_{11 x p}+P_{12 x p} & =P_{11 x h}+P_{12 x h} \\
P_{11 y p}+P_{12 y p} & =P_{11 y h}+P_{12 y h}
\end{aligned}
$$

or:

$$
\left(M_{\mathrm{o}} \cdot \gamma_{v 11 x \mathrm{p}} \cdot v_{11 x \mathrm{p}}\right)+\left(M_{\mathrm{o}} \cdot \gamma_{v 12 x \mathrm{p}} \cdot v_{12 x \mathrm{p}}\right)=
$$

$$
\begin{array}{r}
=\left(M_{\mathrm{o}} \cdot \gamma_{v 11 h} \cdot v_{11 \mathrm{x} h}\right)+\left(M_{\mathrm{o}} \cdot \gamma_{v 12 h} \cdot v_{12 \mathrm{x} h}\right) \\
0 \tag{125}
\end{array}=\left(M_{\mathrm{o}} \cdot \gamma_{v 11 h} \cdot v_{11 \mathrm{yh}}\right)+\left(M_{\mathrm{o}} \cdot \gamma_{v 12 h} \cdot v_{12 \mathrm{yh}}\right) .
$$

Having obtained equation (124) and (125), can determine the conditions of implementation of the law of conservation of momentum for example 2 in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ for the case when the values of the coefficient of proportionality $\gamma_{V}$ (also $\gamma_{v}$ ) lie in the range $\gamma_{V} \geq 1$ (and $\gamma_{v} \geq 1$ ).

Equations (124) and (125) taking into account formula (33) take the form:

$$
\begin{align*}
& \frac{M_{\mathrm{o}} \cdot v_{11 x \mathrm{p}}}{\sqrt{1-\frac{v_{11 \times \mathrm{p}}^{2}}{c_{1}^{2}}}}+\frac{M_{\mathrm{o}} \cdot v_{12 x \mathrm{p}}}{\sqrt{1-\frac{v_{12 x \mathrm{p}}^{2}}{c_{1}^{2}}}}=\frac{M_{\mathrm{o}} \cdot v_{11 \mathrm{x} h}}{\sqrt{1-\frac{v_{11 \mathrm{xh}}^{2}+v_{11 \mathrm{yh}}^{2}}{c_{1}^{2}}}}+\frac{M_{\mathrm{o}} \cdot v_{12 \mathrm{xh}}}{\sqrt{1-\frac{v_{12 \mathrm{xh}}^{2}+v_{12 y h}^{2}}{c_{1}^{2}}}}  \tag{126}\\
& 0=\frac{M_{\mathrm{o}} \cdot v_{11 \mathrm{y} h}}{\sqrt{1-\frac{v_{11 \mathrm{x} h}^{2}+v_{11 \mathrm{yh}}^{2}}{c_{1}^{2}}}}+\frac{M_{\mathrm{o}} \cdot v_{12 \mathrm{y} h}}{\sqrt{1-\frac{v_{12 \mathrm{x} h}^{2}+v_{12 \mathrm{yh}}^{2}}{c_{1}^{2}}}} \tag{127}
\end{align*}
$$

Formulas (103) - (106) and (113) - (116) using the formula (33) can be written:

$$
\begin{gather*}
v_{11 x p}=\frac{V-v_{R}}{1-\frac{V \cdot v_{R}}{c_{1}^{2}}}  \tag{128}\\
v_{12 x p}=\frac{V+v_{R}}{1+\frac{V \cdot v_{R}}{c_{1}^{2}}}  \tag{129}\\
v_{11 x h}=V  \tag{113}\\
v_{11 y h}=-\left(v_{R} \cdot \sqrt{1-\frac{V^{2}}{c_{1}^{2}}}\right)  \tag{130}\\
v_{12 x h}=  \tag{131}\\
1+\frac{V+v_{22 x h}}{c_{1}^{2}}  \tag{132}\\
v_{12 y h}= \\
v_{22 y h} \cdot \sqrt{1-\frac{V^{2}}{c_{1}^{2}}} \\
1+\frac{V \cdot v_{22 x h}}{c_{1}^{2}}
\end{gather*}
$$

By inserting the projections $v_{11 x p}, v_{12 x p}, v_{11 x h}, v_{11 y h}, v_{12 x h}$ and $v_{12 y h}$ of speeds of formulas (113), (128) - (132) in equations (126) and (127) and using the formula (117), we obtain:

$$
\begin{gather*}
\frac{M_{\mathrm{o}} \cdot\left(V-v_{R}\right)}{\sqrt{1-\frac{v_{R}^{2}}{c_{1}^{2}}} \cdot \sqrt{1-\frac{V^{2}}{c_{1}^{2}}}}+\frac{M_{\mathrm{o}} \cdot\left(V+v_{R}\right)}{\sqrt{1-\frac{v_{R}^{2}}{c_{1}^{2}} \cdot \sqrt{1-\frac{V^{2}}{c_{1}^{2}}}}}= \\
=\frac{M_{\mathrm{o}} \cdot V}{\sqrt{1-\frac{v_{R}^{2}}{c_{1}^{2}}} \cdot \sqrt{1-\frac{V^{2}}{c_{1}^{2}}}}+\frac{M_{\mathrm{o}} \cdot\left(V+v_{22 x h}\right)}{\sqrt{1-\frac{v_{R}^{2}}{c_{1}^{2}}} \cdot \sqrt{1-\frac{V^{2}}{c_{1}^{2}}}}  \tag{133}\\
0=-\frac{M_{\mathrm{o}} \cdot v_{R}}{\sqrt{1-\frac{v_{R}^{2}}{c_{1}^{2}}}}+\frac{M_{\mathrm{o}} \cdot v_{22 y h}}{\sqrt{1-\frac{v_{R}^{2}}{c_{1}^{2}}}} \tag{134}
\end{gather*}
$$

or:

$$
\begin{gather*}
V-v_{R}+V+v_{R}=V+V+v_{22 x h}  \tag{135}\\
0=-v_{R}+v_{22 y h} \tag{136}
\end{gather*}
$$

From equations (135) and (136) obtain the necessary conditions (the values of the projections $v_{22 x h}$ and $v_{22 y h}$ of speeds), which in the example 2 will be implemented by law of conservation of momentum in the stationary inertial reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ for the case when the values of the coefficient of proportionality $\gamma_{V}$ are in the range $\gamma_{V} \geq 1$ :

$$
\begin{gather*}
v_{22 x h}=0  \tag{137}\\
v_{22 y h}=v_{R} \tag{138}
\end{gather*}
$$

From (137) and (138) it follows, that the values of projections $v_{22 x h}$ and $v_{22 y h}$ of speeds do not depend on the magnitude of the speed $V$ (and, consequently, do not depend on the magnitude of the coefficient of proportionality $\gamma_{V}$ ).

Substituting conditions (137) and (138) in equations (80) and (81), we obtain:

$$
\begin{equation*}
t_{22 h}=t_{21 h}=0 \tag{139}
\end{equation*}
$$

And substituting equation (139) in the formula (98):

$$
\begin{equation*}
\omega \cdot 0=\left(1-\frac{1}{\gamma_{V}^{2}}\right) \cdot[1+1] \cdot \frac{v_{R}}{V} \tag{140}
\end{equation*}
$$

will have another condition for the implementation of the law of conservation of momentum in the stationary inertial reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ for example 2 :

$$
\begin{equation*}
\gamma_{V}=1 \tag{141}
\end{equation*}
$$

Thus, we can conclude, that in a closed mechanical system of bodies, considered in example 2, for the values of the coefficient of proportionality $\gamma_{V}>1$ the law of conservation of momentum is not satisfied.

Similarly, using equation (124) and (125), can determine the conditions of implementation of the law of conservation of momentum for example 2 in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ for the case, when the values of the coefficient of proportionality $\gamma_{V}$ (also $\gamma_{v}$ ) lie in the range $0<\gamma_{V} \leq 1$ (and $0<\gamma_{v} \leq 1$ ).

Equations (124) and (125) taking into account formula (34) take the form:

$$
\begin{gather*}
\frac{M_{\mathrm{o}} \cdot v_{11 \times \mathrm{p}}}{\sqrt{1+\frac{v_{11 \times \mathrm{p}}^{2}}{c_{2}^{2}}}}+\frac{M_{\mathrm{o}} \cdot v_{12 x \mathrm{p}}}{\sqrt{1+\frac{v_{12 x \mathrm{p}}^{2}}{c_{2}^{2}}}=\frac{M_{\mathrm{o}} \cdot v_{11 \mathrm{x} h}}{\sqrt{1+\frac{v_{11 \mathrm{xh}}^{2}+v_{11 y h}^{2}}{c_{2}^{2}}}}+\frac{M_{\mathrm{o}} \cdot v_{12 \mathrm{xh}}}{\sqrt{1+\frac{v_{12 \mathrm{xh}}^{2}+v_{12 y h}^{2}}{c_{2}^{2}}}}}  \tag{142}\\
0=\frac{M_{\mathrm{o}} \cdot v_{11 \mathrm{yh}}}{\sqrt{1+\frac{v_{11 \mathrm{xh}}^{2}+v_{11 \mathrm{yh}}^{2}}{c_{2}^{2}}}+\frac{M_{\mathrm{o}} \cdot v_{12 \mathrm{xh}}}{\sqrt{1+\frac{v_{12 \mathrm{xh}}^{2}+v_{12 \mathrm{yh}}^{2}}{c_{2}^{2}}}}} \tag{143}
\end{gather*}
$$

Formulas (103) - (106) and (113) - (116) using the formula (34) can be written:

$$
\begin{gather*}
v_{11 x p}=\frac{V-v_{R}}{1+\frac{V \cdot v_{R}}{c_{2}^{2}}}  \tag{144}\\
v_{12 x p}=\frac{V+v_{R}}{1-\frac{V \cdot v_{R}}{c_{2}^{2}}}  \tag{145}\\
v_{11 \times h}=V \tag{113}
\end{gather*}
$$

$$
\begin{gather*}
v_{11 y h}=-\left(v_{R} \cdot \sqrt{1+\frac{V^{2}}{c_{2}^{2}}}\right)  \tag{146}\\
v_{12 x h}=\frac{V+v_{22 x h}}{1-\frac{V \cdot v_{22 x h}}{c_{2}^{2}}}  \tag{147}\\
v_{12 y h}=\frac{v_{22 y h} \cdot \sqrt{1+\frac{V^{2}}{c_{2}^{2}}}}{1-\frac{V \cdot v_{22 x h}}{c_{2}^{2}}} \tag{148}
\end{gather*}
$$

By inserting the projections $v_{11 x p}, v_{12 x p}, v_{11 x h}, v_{11 y h}, v_{12 x h}$ and $v_{12 y h}$ of speeds of formulas (113), (144)-(148) in equations (142) and (143) and using the formula (117), we obtain:

$$
\begin{gather*}
\frac{M_{\mathrm{o}} \cdot\left(V-v_{R}\right)}{\sqrt{1+\frac{v_{R}^{2}}{c_{2}^{2}}} \cdot \sqrt{1+\frac{V^{2}}{c_{2}^{2}}}}+\frac{M_{\mathrm{o}} \cdot\left(V+v_{R}\right)}{\sqrt{1+\frac{v_{R}^{2}}{c_{2}^{2}} \cdot \sqrt{1+\frac{V^{2}}{c_{2}^{2}}}}}= \\
=\frac{M_{\mathrm{o}} \cdot V}{\sqrt{1+\frac{v_{R}^{2}}{c_{2}^{2}}} \cdot \sqrt{1+\frac{V^{2}}{c_{2}^{2}}}}+\frac{M_{\mathrm{o}} \cdot\left(V+v_{22 x h}\right)}{\sqrt{1+\frac{v_{R}^{2}}{c_{2}^{2}}} \cdot \sqrt{1+\frac{V^{2}}{c_{2}^{2}}}}  \tag{149}\\
0=-\frac{M_{\mathrm{o}} \cdot v_{R}}{\sqrt{1+\frac{v_{R}^{2}}{c_{2}^{2}}}}+\frac{M_{\mathrm{o}} \cdot v_{22 y h}}{\sqrt{1+\frac{v_{R}^{2}}{c_{2}^{2}}}} \tag{150}
\end{gather*}
$$

or:

$$
\begin{gather*}
V-v_{R}+V+v_{R}=V+V+v_{22 x h}  \tag{135}\\
0=-v_{R}+v_{22 y h} \tag{136}
\end{gather*}
$$

From equations (135) and (136) obtain the necessary conditions (the values of the projections $v_{22 x h}$ and $v_{22 y h}$ of speeds), which in the example 2 will be implemented by law of conservation of momentum in the stationary inertial reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ for the case when the values of the coefficient of proportionality $\gamma_{V}$ are in the range $0<\gamma_{V} \leq 1$ :

$$
\begin{gather*}
v_{22 x h}=0  \tag{137}\\
v_{22 y h}=v_{R} \tag{138}
\end{gather*}
$$

Equalities (137) and (138), as has been shown in considering the case, when the value of the coefficient of proportionality $\gamma_{V}$ are in the range $\gamma_{V} \geq 1$, lead to the same condition for the implementation of the law of conservation of momentum for example 2 in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ for the case, when the values of the coefficient of proportionality $\gamma_{V}$ are in the range $0<\gamma_{V} \leq 1:$

$$
\begin{equation*}
\gamma_{V}=1 \tag{141}
\end{equation*}
$$

Consequently, we can make a similar conclusion, that in a closed mechanical system of bodies, considered in example 2, for the values of the coefficient of proportionality $0<\gamma_{V}<1$ the law of conservation of momentum also is not satisfied.

Summarizing the findings, we note, that in a closed mechanical system of bodies, considered in example 2, for the values of the coefficient of proportionality $\gamma_{V}>1$ and $0<\gamma_{V}<1$ the law of conservation of momentum is not satisfied.

The law of conservation of momentum will be carried out only, if the coefficient of proportionality $\gamma_{V}$ equal to 1 .

In the case of the obligation to fulfill the law of conservation of momentum of a closed mechanical system of bodies, considered in example 2, based on the formulas (26) - (28), and given, that the coefficient of proportionality $\gamma_{V}=1$, constants $c_{1}$ and $c_{2}$ will have the following meanings:

$$
\begin{align*}
& c_{1}= \pm \infty  \tag{151}\\
& c_{2}= \pm \infty \tag{152}
\end{align*}
$$

## 8. Score quantities of momentums in example 3

Given the possible observation, that in example 2 the failure of the law of conservation of momentum can be attributed to the assumptions of infinitesimal mass of the string 3 , consider the example 3 .

Example 3 differs from example 2 in that in example 3 the mass of the string 3 is not infinitely small.

For example 3 will try to assess the impact of magnitude of the momentum of the string 3 on the magnitude of the momentum of bodies 1 and 2 and string 3.

Assume that there are two inertial reference systems, similar to those of reference systems, shown in Fig.1, stationary $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and mobile $\mathrm{O}_{2} x_{2} y_{2} z_{2}$, which moves with speed $V$ parallel to the axis $\mathrm{O}_{1} x_{1}$ relative to the system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$.

Suppose that there is a closed mechanical system of bodies, shown in Fig. 13 and consisting of point bodies 1 and 2, with equal mass $M_{\mathrm{o}}$ at rest, and a string 3 .

Bodies 1 and 2 are connected by a string 3, which has a mass of uniformly distributed along its length and equal to $m_{0}$ at rest.

Bodies 1 and 2 rotate with angular speed $\omega$ around a common center of mass - the point O .

Distance from the point body 1 (body 2$)$ to point O is equal to $R$.
Let's put a closed mechanical system of bodies 1 and 2 with a string 3 in the moving reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ so, that the point O would be stationary in this reference system, and coincided with the origin $\mathrm{O}_{2}$, and the rotation of bodies 1 and 2 around it would occur in a clockwise direction in the plane of $\mathrm{O}_{2} x_{2} y_{2}$, as shown in Fig. 14.

Also assume, that at the start of timing $\left(t_{2}=0\right)$ in the reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ bodies 1 and 2 were on the axis $\mathrm{O}_{2} x_{2}$, with the body 1 had a positive coordinate, and the body 2 - negative.

In the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at any time $t_{2}$ bodies 1 and 2 will have the speeds $v_{21}$ and $v_{22}$, equal $v_{R}$ :

$$
\begin{equation*}
v_{21}=v_{22}=v_{R}=\omega \cdot R \tag{77}
\end{equation*}
$$

In this case, the projections $v_{21 x}$ and $v_{21 y}$ of speed of the body 1 and the projections $v_{22 x}$ and $v_{22 y}$ of speed of body 2 on the axis $\mathrm{O}_{2} x_{2}$ and $\mathrm{O}_{2} y_{2}$,
respectively, for time $t_{2}$ will be equal to:

$$
\begin{gather*}
v_{21 x}=-\left[v_{R} \cdot \sin \left(\omega \cdot t_{2}\right)\right]  \tag{153}\\
v_{21 y}=-\left[v_{R} \cdot \cos \left(\omega \cdot t_{2}\right)\right]  \tag{154}\\
v_{22 x}=v_{R} \cdot \sin \left(\omega \cdot t_{2}\right)  \tag{155}\\
v_{22 y}=v_{R} \cdot \cos \left(\omega \cdot t_{2}\right) \tag{156}
\end{gather*}
$$

The relationship between the coordinates $x_{21}$ and $y_{21}$ of the body 1 and the relationship between the coordinates $x_{22}$ and $y_{22}$ body 2 depending on the time $t_{2}$ in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ can be written as:

$$
\begin{align*}
x_{21} & =R \cdot \cos \left(\omega \cdot t_{2}\right)  \tag{157}\\
y_{21} & =-\left[R \cdot \sin \left(\omega \cdot t_{2}\right)\right]  \tag{158}\\
x_{22} & =-\left[R \cdot \cos \left(\omega \cdot t_{2}\right)\right]  \tag{159}\\
y_{22} & =R \cdot \sin \left(\omega \cdot t_{2}\right) \tag{160}
\end{align*}
$$

Similarly, for the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ you can get the dependencies:

- dependencies of the projections $v_{21 \text { xpi }}$ and $v_{21 \text { ypi }}$ of the speed of the $i$-point of the string 3 , which is located at a distance $\rho_{i}$ from point O on the segment from point O to the body 1 , on the axis $\mathrm{O}_{2} x_{2}$ and $\mathrm{O}_{2} y_{2}$ on the time $t_{2}$ :

$$
\begin{align*}
& v_{21 \times \rho i}=-\left[v_{R} \cdot \frac{\rho_{i}}{R} \cdot \sin \left(\omega \cdot t_{2}\right)\right]  \tag{161}\\
& v_{21 y \rho i}=-\left[v_{R} \cdot \frac{\rho_{i}}{R} \cdot \cos \left(\omega \cdot t_{2}\right)\right] \tag{162}
\end{align*}
$$

- dependencies of the projections $v_{22 x p j}$ and $v_{22 y p j}$ of the speed of the $j$-point of the string 3 , which is located at a distance $\rho_{j}$ from point O on the segment from point O to the body 2 , on the axis $\mathrm{O}_{2} x_{2}$ and $\mathrm{O}_{2} y_{2}$ on the time $t_{2}$ :

$$
\begin{align*}
& v_{22 x \rho j}=v_{R} \cdot \frac{\rho_{j}}{R} \cdot \sin \left(\omega \cdot t_{2}\right)  \tag{163}\\
& v_{22 y \rho j}=v_{R} \cdot \frac{\rho_{j}}{R} \cdot \cos \left(\omega \cdot t_{2}\right) \tag{164}
\end{align*}
$$

- dependencies of the values of the coordinates $x_{21 \rho i}$ and $y_{21 \rho i}$ of $i$-point of the string 3 and the coordinates $x_{22 \rho j}$ and $y_{22 \rho j}$ of $j$-point of the string 3 on the time $t_{2}$ :

$$
\begin{gather*}
x_{21 \rho i}=\rho_{i} \cdot \cos \left(\omega \cdot t_{2}\right)  \tag{165}\\
y_{21 \rho i}=-\left[\rho_{i} \cdot \sin \left(\omega \cdot t_{2}\right)\right]  \tag{166}\\
x_{22 \rho j}=-\left[\rho_{j} \cdot \cos \left(\omega \cdot t_{2}\right)\right]  \tag{167}\\
y_{22 \rho j}=\rho_{j} \cdot \sin \left(\omega \cdot t_{2}\right) \tag{168}
\end{gather*}
$$

Now you can proceed to consider the movement of bodies 1 and 2 and string 3 in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$.

Suppose that, as shown in Fig.17, the mobile inertial reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ moves with speed $V$ relative to the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$, where in the two systems as the origin of time ( $t_{1}=0$ and $t_{2}=0$ ) is selected such a time, when origins $\mathrm{O}_{1}$ and $\mathrm{O}_{2}$ of these systems are the same (ie, when points $\mathrm{O}_{1}$, $\mathrm{O}_{2}$ and O are the same).


Fig. 17

From equations (1)-(3), (5)-(7), (9), to consider the motion of the body 1 , we can write the following:

- relationships between coordinates $x_{11}$ and $y_{11}$ of the body 1 at time $t_{1}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and coordinates $x_{21}$ and $y_{21}$ of the body 1 in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{2}$, which corresponds to the time $t_{1}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ :

$$
\begin{align*}
x_{11} & =\gamma_{V} \cdot\left[x_{21}+\left(V \cdot t_{2}\right)\right]  \tag{169}\\
x_{21} & =\gamma_{V} \cdot\left[x_{11}-\left(V \cdot t_{1}\right)\right] \tag{170}
\end{align*}
$$

$$
\begin{equation*}
y_{11}=y_{21} \tag{87}
\end{equation*}
$$

- relationship between the values of times $t_{1}$ and $t_{2}$ in describing the motion of the body 1 :

$$
\begin{align*}
t_{1} & =\frac{\left(\gamma_{V}^{2}-1\right) \cdot x_{21}}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{2}\right)  \tag{171}\\
t_{2} & =\frac{\left(1-\gamma_{V}^{2}\right) \cdot x_{11}}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{1}\right) \tag{172}
\end{align*}
$$

while taking into account the equation (157) formula (171) becomes:

$$
\begin{equation*}
t_{1}=\frac{\left(\gamma_{V}^{2}-1\right) \cdot R \cdot \cos \left(\omega \cdot t_{2}\right)}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{2}\right) \tag{173}
\end{equation*}
$$

- relationships between the projections $v_{x 11}$ and $v_{y 11}$ of speed $v_{11}$ of the body 1 at time $t_{1}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and similar projections $v_{x 21}$ and $v_{y 21}$ of speed $v_{21}$ of the body 1 in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{2}$, which corresponds to the time $t_{1}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ :

$$
\begin{align*}
& v_{x 11}=\frac{v_{x 21}+V}{\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{x 21}}{\gamma_{V}^{2} \cdot V}+1}  \tag{174}\\
& v_{y 11}=\frac{v_{y 21}}{\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{x 21}}{\gamma_{V} \cdot V}+\gamma_{V}} \tag{175}
\end{align*}
$$

Similarly, for the consideration of motion of the body 2 can be written:

- relationships between coordinates $x_{12}$ and $y_{12}$ of the body 2 at time $t_{1}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and coordinates $x_{22}$ and $y_{22}$ of the body 2 in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{2}$, which corresponds to the time $t_{1}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ :

$$
\begin{gather*}
x_{12}=\gamma_{V} \cdot\left[x_{22}+\left(V \cdot t_{2}\right)\right]  \tag{176}\\
x_{22}=\gamma_{V} \cdot\left[x_{12}-\left(V \cdot t_{1}\right)\right]  \tag{177}\\
y_{12}=y_{22} \tag{89}
\end{gather*}
$$

- relationship between the values of times $t_{1}$ and $t_{2}$ in describing the motion of the body 2 :

$$
\begin{align*}
t_{1} & =\frac{\left(\gamma_{V}^{2}-1\right) \cdot x_{22}}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{2}\right)  \tag{178}\\
t_{2} & =\frac{\left(1-\gamma_{V}^{2}\right) \cdot x_{12}}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{1}\right) \tag{179}
\end{align*}
$$

while taking into account the equation (159) formula (178) becomes:

$$
\begin{equation*}
t_{1}=-\frac{\left(\gamma_{V}^{2}-1\right) \cdot R \cdot \cos \left(\omega \cdot t_{2}\right)}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{2}\right) \tag{180}
\end{equation*}
$$

- relationships between the projections $v_{x 12}$ and $v_{y 12}$ of speed $v_{12}$ of the body 2 at time $t_{1}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and similar projections $v_{x 22}$ and $v_{y 22}$ of speed $v_{22}$ of the body 2 in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{2}$, which corresponds to the time $t_{1}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ :

$$
\begin{align*}
v_{x 12} & =\frac{v_{x 22}+V}{\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{x 22}}{\gamma_{V}^{2} \cdot V}+1}  \tag{181}\\
v_{y 12} & =\frac{v_{y 22}}{\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{x 22}}{\gamma_{V} \cdot V}+\gamma_{V}} \tag{182}
\end{align*}
$$

Also to consider the motion of $i$-the point of the string 3 , which is located at a distance $\rho_{i}$ from point O on the segment from point O to the body 1 in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$, we can write the following:

- relationships between coordinates $x_{11 p i}$ and $y_{11 p i}$ of $i$-the point of the string 3 at time $t_{1}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and coordinates $x_{21 \rho i}$ and $y_{21 \rho i}$ of $i$-the point of the string 3 in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{2}$, which corresponds to the time $t_{1}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ :

$$
\begin{align*}
x_{11 \rho i} & =\gamma_{V} \cdot\left[x_{21 \rho i}+\left(V \cdot t_{2}\right)\right]  \tag{183}\\
x_{21 \rho i} & =\gamma_{V} \cdot\left[x_{11 \rho i}-\left(V \cdot t_{1}\right)\right]  \tag{184}\\
y_{11 \rho i} & =y_{21 \rho i} \tag{185}
\end{align*}
$$

- relationship between the values of times $t_{1}$ and $t_{2}$ in describing the motion of $i$-the point of the string 3 :

$$
\begin{align*}
t_{1} & =\frac{\left(\gamma_{V}^{2}-1\right) \cdot x_{21 \rho i}}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{2}\right)  \tag{186}\\
t_{2} & =\frac{\left(1-\gamma_{V}^{2}\right) \cdot x_{11 \rho i}}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{1}\right) \tag{187}
\end{align*}
$$

while taking into account the equation (165) formula (186) becomes:

$$
t_{1}=\frac{\left(\gamma_{V}^{2}-1\right) \cdot \rho_{i} \cdot \cos \left(\omega \cdot t_{2}\right)}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{2}\right)
$$

- relationships between the projections $v_{x 11 \rho i}$ and $v_{y 11 \rho i}$ of speed $v_{11 \rho i}$ of $i$-the point of the string 3 at time $t_{1}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and similar projections $v_{x 21 \rho i}$ and $v_{y 21 \rho i}$ of speed $v_{21 \rho i}$ of $i$-the point of the string 3 in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{2}$, which corresponds to the time $t_{1}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ :

$$
\begin{align*}
& v_{x 11 \rho i}=\frac{v_{x 21 \rho i}+V}{\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{x 21 \rho i}}{\gamma_{V}^{2} \cdot V}+1}  \tag{189}\\
& v_{y 11 \rho i}=\frac{v_{y 21 \rho i}}{\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{x 21 \rho i}}{\gamma_{V} \cdot V}+\gamma_{V}} \tag{190}
\end{align*}
$$

Also to consider the motion of $j$-the point of the string 3 , which is located at a distance $\rho_{j}$ from point O on the segment from point O to the body 2 in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$, we can write the following:

- relationships between coordinates $x_{12 \rho j}$ and $y_{12 \rho j}$ of $j$-the point of the string 3 at time $t_{1}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and coordinates $x_{22 \rho j}$ and $y_{22 \rho j}$ of $j$-the point of the string 3 in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{2}$, which corresponds to the time $t_{1}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ :

$$
\begin{align*}
& x_{12 \rho j}=\gamma_{V} \cdot\left[x_{22 \rho j}+\left(V \cdot t_{2}\right)\right]  \tag{191}\\
& x_{22 \rho j}=\gamma_{V} \cdot\left[x_{12 \rho j}-\left(V \cdot t_{1}\right)\right]  \tag{192}\\
& y_{12 \rho j}=y_{22 \rho j}
\end{align*}
$$

- relationship between the values of times $t_{1}$ and $t_{2}$ in describing the motion of $j$-the point of the string 3 :

$$
\begin{align*}
t_{1} & =\frac{\left(\gamma_{V}^{2}-1\right) \cdot x_{22 \rho j}}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{2}\right)  \tag{194}\\
t_{2} & =\frac{\left(1-\gamma_{V}^{2}\right) \cdot x_{12 \rho j}}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{1}\right) \tag{195}
\end{align*}
$$

while taking into account the equation (167) formula (194) becomes:

$$
\begin{equation*}
t_{1}=-\frac{\left(\gamma_{V}^{2}-1\right) \cdot \rho_{j} \cdot \cos \left(\omega \cdot t_{2}\right)}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{2}\right) \tag{196}
\end{equation*}
$$

- relationships between the projections $v_{x 12 \rho j}$ and $v_{y 12 \rho j}$ of speed $v_{12 \rho j}$ of $j$-the point of the string 3 at time $t_{1}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ and similar projections $v_{x 22 \rho j}$ and $v_{y 22 \rho j}$ of speed $v_{22 \rho j}$ of $j$-the point of the string 3 in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{2}$, which corresponds to the time $t_{1}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ :

$$
\begin{align*}
& v_{x 12 \rho j}=\frac{v_{x 22 \rho j}+V}{\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{x 22 \rho j}}{\gamma_{V}^{2} \cdot V}+1}  \tag{197}\\
& v_{y 12 \rho j}=\frac{v_{y 22 \rho j}}{\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{x 22 \rho j}}{\gamma_{V} \cdot V}+\gamma_{V}} \tag{198}
\end{align*}
$$

In order to initiate testing of the law of conservation of momentum must select two points in time in the stationary inertial reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$.

First time - this is $t_{1 p}$.
Suppose that in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ at time $t_{1}$, equal to $t_{1 p}$, the bodies 1 and 2 are on the line parallel to the axis $\mathrm{O}_{1} y_{1}$ (or coinciding with it), ie where:

$$
\begin{equation*}
x_{11}=x_{12} \tag{199}
\end{equation*}
$$

Condition (199) is possible only in the case, when in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at the time $t_{2}$, equal to $t_{2 p}$, corresponding to the time $t_{1 p}$ in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$, the following conditions:

$$
\begin{align*}
& x_{21}=x_{22}  \tag{200}\\
& \omega \cdot t_{2 p}=\frac{\pi}{2} \tag{201}
\end{align*}
$$

As shown in Fig.15, according to equations (201), (153) - (156) in the
mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{2 p}$ the bodies 1 and 2 , respectively, have the following values of the projections $v_{21 x p}, v_{21 y p}$ and $v_{22 x p}, v_{22 y p}$ of speeds of his movement on the axis $\mathrm{O}_{2} x_{2}$ and $\mathrm{O}_{2} y_{2}$ :

$$
\begin{align*}
& v_{21 \times p}=-v_{R}  \tag{99}\\
& v_{21 y p}=0  \tag{100}\\
& v_{22 x p}=v_{R}  \tag{101}\\
& v_{22 y p}=0 \tag{102}
\end{align*}
$$

And in accordance with equations (201), (161)-(164) in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{2 p}$ the $i$-the point of the string 3 , which is located at a distance $\rho_{i}$ from point O on the segment from point O to the body 1 , and the $j$-the point of the string 3 , which is located at a distance $\rho_{j}$ from point O on the segment from point O to the body 2 , respectively, have the following values of the projections $v_{21 \text { xpip }}, v_{21 \text { ypip }}$ and $v_{22 \text { xpjp }}, v_{22 \text { ypjp }}$ of speeds of his movement on the axis $\mathrm{O}_{2} x_{2}$ and $\mathrm{O}_{2} y_{2}$ :

$$
\begin{array}{r}
v_{21 \times \rho i \mathrm{p}}=-\left(v_{R} \cdot \frac{\rho_{i}}{R}\right) \\
v_{21 y \rho i \mathrm{p}}=0 \\
v_{22 x \rho j \mathrm{p}}=v_{R} \cdot \frac{\rho_{j}}{R} \\
v_{22 y \rho j \mathrm{p}}=0 \tag{205}
\end{array}
$$

A second point in time we choose $t_{1 h}$.
Suppose that, as shown in Fig.16, in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ at time $t_{1}$, equal to $t_{1 h}$, the position of body 1 will be consistent the position of the body 1 at time $t_{2}$, equal to $t_{21 h}$ :

$$
\begin{equation*}
t_{21 h}=0 \tag{206}
\end{equation*}
$$

in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$, ie when the body 1 will be on the axis $\mathrm{O}_{2} x_{2}$.

The value of time $t_{1 h}$ can be determined from equation (173) on the basis of conditions (206):

$$
\begin{equation*}
t_{1 h}=\frac{\left(\gamma_{V}^{2}-1\right) \cdot R}{\gamma_{V} \cdot V} \tag{207}
\end{equation*}
$$

According to equations (206), (153), (154) in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{21 h}$ body 1 will have the following values of the projections $v_{21 \times h}$ and $v_{21 y h}$ of the speed of his movement on the axis $\mathrm{O}_{2} x_{2}$ and $\mathrm{O}_{2} y_{2}$ :

$$
\begin{gather*}
v_{21 x h}=0  \tag{111}\\
v_{21 y h}=-v_{R} \tag{112}
\end{gather*}
$$

In the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ at the time $t_{1}$, equal to $t_{1 h}$, the position of body 2 will be consistent the position of the body 2 in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{2}$, equal to $t_{22 h}$, which can be determined on the basis of equations (180) and (207):

$$
\begin{equation*}
\frac{\left(\gamma_{V}^{2}-1\right) \cdot R}{\gamma_{V} \cdot V}=-\frac{\left(\gamma_{V}^{2}-1\right) \cdot R \cdot \cos \left(\omega \cdot t_{22 h}\right)}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{22 h}\right) \tag{208}
\end{equation*}
$$

or:

$$
\begin{equation*}
\left(\omega \cdot t_{22 h}\right)=\frac{\left(\gamma_{V}^{2}-1\right) \cdot\left[1+\cos \left(\omega \cdot t_{22 h}\right)\right] \cdot v_{R}}{\gamma_{V}^{2} \cdot V} \tag{209}
\end{equation*}
$$

Similar in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ at the time $t_{1}$, equal to $t_{1 h}$, the position of the $i$-the point of the string 3 will be consistent the position of the $i$-the point of the string 3 , which is located at a distance $\rho_{i}$ from point O on the segment from point O to the body 1 , in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{2}$, equal to $t_{21 \rho i h}$, which can be determined on the basis of equations (188) and (207):

$$
\begin{equation*}
\frac{\left(\gamma_{V}^{2}-1\right) \cdot R}{\gamma_{V} \cdot V}=\frac{\left(\gamma_{V}^{2}-1\right) \cdot \rho_{i} \cdot \cos \left(\omega \cdot t_{21 \rho i h}\right)}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{21 \rho i h}\right) \tag{210}
\end{equation*}
$$

or:

$$
\begin{equation*}
\left(\omega \cdot t_{21 \rho i h}\right)=\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{R}}{\gamma_{V} \cdot V} \cdot\left\{1-\left[\frac{\rho_{\mathrm{i}}}{\mathrm{R}} \cdot \cos \left(\omega \cdot t_{21 \rho i h}\right)\right]\right\} \tag{211}
\end{equation*}
$$

Also in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ at the time $t_{1}$, equal to $t_{1 h}$, the position of the $j$-the point of the string 3 will be consistent the position of the $j$-the point of the string 3 , which is located at a distance $\rho_{j}$ from point O on
the segment from point O to the body 2 , in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ at time $t_{2}$, equal to $t_{22 \rho j h}$, which can be determined on the basis of equations (196) and (207):

$$
\begin{equation*}
\frac{\left(\gamma_{V}^{2}-1\right) \cdot R}{\gamma_{V} \cdot V}=-\frac{\left(\gamma_{V}^{2}-1\right) \cdot \rho_{j} \cdot \cos \left(\omega \cdot t_{22 \rho j h}\right)}{\gamma_{V} \cdot V}+\left(\gamma_{V} \cdot t_{22 \rho j h}\right) \tag{212}
\end{equation*}
$$

or:

$$
\begin{equation*}
\left(\omega \cdot t_{22 \rho j h}\right)=\frac{\left(\gamma_{V}^{2}-1\right) \cdot v_{R}}{\gamma_{V}^{2} \cdot V} \cdot\left\{1+\left[\frac{\rho_{\mathrm{j}}}{\mathrm{R}} \cdot \cos \left(\omega \cdot t_{22 \rho j h}\right)\right]\right\} \tag{213}
\end{equation*}
$$

To handle complex calculations in equations (209), (211) and (213), the values of momentums will try to determine by simple numerical examples.

For consideration in the mobile reference system $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ the string 3 conditionally divided by 17 equal parts ( $i=0,1,2,3,4,5,6,7,8$ and $j=1,2,3$, $4,5,6,7,8)$ with accommodation in the center of each part the point of the body with rest mass $m_{017}$, equal to:

$$
\begin{equation*}
m_{017}=\frac{m_{0}}{17} \tag{214}
\end{equation*}
$$

In this case, the distance $\rho_{i}$ from point O to the $i$-the point of the string 3 , located on the segment from point $O$ to the body 1 , will be equal to:

$$
\begin{equation*}
\rho_{i}=\frac{2 \cdot i}{17} \tag{215}
\end{equation*}
$$

And the distance $\rho_{j}$ from point O to the $j$-the point of the string 3 , located on the segment from point O to the body 2 , will be equal to:

$$
\begin{equation*}
\rho_{j}=\frac{2 \cdot j}{17} \tag{216}
\end{equation*}
$$

First, consider the case, when the values of the coefficient of proportionality $\gamma_{V}$ (also $\gamma_{v}$ ) lie in the range $\gamma_{V} \geq 1$ (and $\gamma_{v} \geq 1$ ).

In the case, if the values of the coefficient of proportionality $\gamma_{v}$ are in the range $\gamma_{v} \geq 1$, then as follows from formula (72), in any inertial reference system $O x y z$ the projections $P_{x}$ and $P_{y}$ of the momentum of a material point, moving with the speed $v$ and having a rest mass $m_{0}$, on the axis $O x$ and $O y$, respectively, can be written:

$$
\begin{align*}
P_{x} & =\frac{54}{} \frac{m_{\boldsymbol{o}} v_{x}}{\sqrt{1-\frac{\left(v_{x}^{2}+v_{y}^{2}\right)}{c_{1}^{2}}}} \\
P_{\mathrm{y}} & =\frac{m_{\boldsymbol{o}} v_{y}}{\sqrt{1-\frac{\left(v_{x}^{2}+v_{y}^{2}\right)}{c_{1}^{2}}}} \tag{217}
\end{align*}
$$

where: $v_{x}$ and $v_{y}$ - the projections of the speed $v$ of a material point on the axis $O x$ and $O y$, respectively.

Assume in the considered example 3 (shown in Fig. 13 - Fig.17), that:

$$
\begin{align*}
& \frac{V}{c_{1}}=0,9  \tag{219}\\
& \frac{v_{R}}{c_{1}}=0,8  \tag{220}\\
& \frac{m_{0}}{M_{0}}=0,1 \tag{221}
\end{align*}
$$

To determine the values of the momentums of the system of bodies 1 and 2 and string 3 in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ at time $t_{1 p}$ will use equation (99) - (102), (217), (218), the raw data (214) - (216), (219) - (221) and the formulas, derived from the equations (174), (175), (181), (182), (189), (190), (197) and (198), taking into account equation (33) :

$$
\begin{align*}
& v_{x 11}= \frac{v_{x 21}+V}{1+\frac{V \cdot v_{x 21}}{c_{1}^{2}}}  \tag{222}\\
& v_{y 11}= v_{y 21} \cdot \sqrt{1-\frac{V^{2}}{c_{1}^{2}}}  \tag{223}\\
& 1+\frac{V \cdot v_{x 21}}{c_{1}^{2}}
\end{align*}, \begin{gathered}
v_{x 22}+V  \tag{224}\\
v_{x 12}+\frac{V \cdot v_{x 22}}{c_{1}^{2}}
\end{gathered}
$$

$$
\begin{align*}
& v_{y 12}=\frac{v_{y 22} \cdot \sqrt{1-\frac{V^{2}}{c_{1}^{2}}}}{1+\frac{V \cdot v_{x 22}}{c_{1}^{2}}}  \tag{225}\\
& v_{x 11 \rho i}=\frac{v_{x 21 \rho i}+V}{1+\frac{V \cdot v_{x 21 \rho i}}{c_{1}^{2}}}  \tag{226}\\
& v_{y 11 \rho i}=\frac{v_{y 21 \rho i} \cdot \sqrt{1-\frac{V^{2}}{c_{1}^{2}}}}{1+\frac{V \cdot v_{x 21 \rho i}}{c_{1}^{2}}}  \tag{227}\\
& v_{x 12 \rho j}=\frac{v_{x 22 \rho j}+V}{1+\frac{V \cdot v_{x 22 \rho j}}{c_{1}^{2}}}  \tag{228}\\
& v_{y 12 \rho j}=\frac{v_{y 22 \rho j} \cdot \sqrt{1-\frac{V^{2}}{c_{1}^{2}}}}{1+\frac{V \cdot v_{x 22 \rho j}}{c_{1}^{2}}} \tag{229}
\end{align*}
$$

The results of digital calculations presented in Tab.7.

Tab. 7
Range $\gamma_{V} \geq 1$ (and $\gamma_{v} \geq 1$ ). Time $t_{1 p}$.

| Object | Mobile reference system <br> $\mathrm{O}_{2} x_{2} y_{2} z_{2}$$\|$The projections of the <br> velocity <br> (the dimension $\mathrm{c}_{1}$ ) |  | Stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | The projections of the velocity (the dimension $\mathrm{c}_{1}$ ) |  | Projections of the momentum <br> (the dimension $\mathrm{c}_{1} \cdot M_{0}$ ) |  |
|  | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{2} x_{2} \\ \hline \end{gathered}$ | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{2} y_{2} \\ \hline \end{gathered}$ | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{1} x_{1} \\ \hline \end{gathered}$ | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{1} y_{1} \\ \hline \end{gathered}$ | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{1} x_{1} \\ \hline \end{gathered}$ | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{1} y_{1} \\ \hline \end{gathered}$ |
| Body 1 | -0,8 | 0 | 0,3571429 | 0 | 0,3823596 | 0 |
| Body 2 | 0,8 | 0 | 0,9883721 | 0 | 6,5001125 | 0 |
| Bodies 1 and 2 |  |  |  |  | 6,8824472 | 0 |
| $i=0$ | 0 | 0 | 0,9 | 0 | 0,0121455 | 0 |
| $i=1$ | -0,09412 | 0 | 0,8804627 | 0 | 0,0109239 | 0 |
| $i=2$ | -0,18824 | 0 | 0,8569405 | 0 | 0,0097801 | 0 |
| $i=3$ | -0,28235 | 0 | 0,8280757 | 0 | 0,0086887 | 0 |
| $i=4$ | -0,37647 | 0 | 0,7918149 | 0 | 0,0076261 | 0 |
| $i=5$ | -0,47059 | 0 | 0,744898 | 0 | 0,0065676 | 0 |
| $i=6$ | -0,56471 | 0 | 0,6818182 | 0 | 0,0054827 | 0 |
| $i=7$ | -0,65882 | 0 | 0,5924855 | 0 | 0,0043263 | 0 |
| $i=8$ | -0,75294 | 0 | 0,4562044 | 0 | 0,0030156 | 0 |
| $j=1$ | 0,094118 | 0 | 0,9164859 | 0 | 0,0134755 | 0 |
| $j=2$ | 0,188235 | 0 | 0,9305835 | 0 | 0,0149531 | 0 |
| $j=3$ | 0,282353 | 0 | 0,99427767 | 0 | 0,0166327 | 0 |
| $j=4$ | 0,376471 | 0 | 0,9534271 | 0 | 0,0185940 | 0 |
| $j=5$ | 0,470588 | 0 | 0,9628099 | 0 | 0,0209623 | 0 |
| $j=6$ | 0,564706 | 0 | 0,9711388 | 0 | 0,0239506 | 0 |
| $j=7$ | 0,658824 | 0 | 0,978582 | 0 | 0,0279629 | 0 |
| $j=8$ | 0,752941 | 0 | 0,9852735 | 0 | 0,0338959 | 0 |
| String 3 |  |  |  |  | 0,2389836 | 0 |
| Bodies 1 and 2 and string 3 |  |  |  |  | 7,1214557 | 0 |

To determine the values of the momentums of the system of bodies 1 and 2 and string 3 in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ at time $t_{1 h}$ will use
equation (155)-(156), (161)-(164), (111)-(112), (217)-(218), (222)-(229), the raw data (219)-(221), (214)-(216) and the formulas, derived from the equations (209), (211) and (213), taking into account equation (33) :

$$
\begin{array}{r}
\left(\omega \cdot t_{22 h}\right)=\frac{V \cdot v_{R} \cdot\left[1+\cos \left(\omega \cdot t_{22 h}\right)\right]}{c_{1}^{2}} \\
\left(\omega \cdot t_{21 \rho i h}\right)=\frac{V \cdot v_{R}}{c_{1}^{2}} \cdot\left\{1-\left[\frac{\rho_{i}}{R} \cdot \cos \left(\omega \cdot t_{21 \rho i h}\right)\right]\right\} \\
\left(\omega \cdot t_{22 \rho j h}\right)=\frac{V \cdot v_{R}}{c_{1}^{2}} \cdot\left\{1+\left[\frac{\rho_{j}}{R} \cdot \cos \left(\omega \cdot t_{22 \rho j h}\right)\right]\right\} \tag{232}
\end{array}
$$

The results of digital calculations presented in Tab.8.

Tab. 8
Range $\gamma_{V} \geq 1$ (and $\gamma_{v} \geq 1$ ). Time $t_{1 h}$.

| Object | Mobile re <br> O | nce system $y_{2} z_{2}$ | Stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | The projections of the velocity <br> (the dimension $\mathrm{c}_{1}$ ) |  | The projections of the velocity <br> (the dimension $\mathrm{c}_{1}$ ) |  | Projections of the momentum <br> (the dimension $\mathrm{c}_{1} \cdot M_{0}$ ) |  |
|  | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{2} x_{2} \\ \hline \end{gathered}$ | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{2} y_{2} \\ \hline \end{gathered}$ | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{1} x_{1} \end{gathered}$ | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{1} y_{1} \\ \hline \end{gathered}$ | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{1} x_{1} \end{gathered}$ | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{1} y_{1} \\ \hline \end{gathered}$ |
| Body 1 | 0 | -0,8 | 0,9 | -0,34871 | 3,441236 | -1,333333 |
| Body 2 | 0,700743 | 0,385953 | 0,9816482 | 0,103168 | 6,1205934 | 0,6432543 |
| Bodies 1 and 2 |  |  |  |  | 9,5618294 | -0,690079 |
| $i=0$ | 0 | 0 | 0,9 | 0 | 0,0121455 | 0 |
| $i=1$ | -0,05716 | -0,07477 | 0,8885503 | -0,03436 | 0,0114249 | -0,000442 |
| $i=2$ | -0,10286 | -0,15765 | 0,8784626 | -0,07573 | 0,0109532 | -0,000944 |
| $i=3$ | -0,13452 | -0,24825 | 0,8709212 | -0,12311 | 0,0107684 | -0,001522 |
| $i=4$ | -0,14977 | -0,3454 | 0,8671108 | -0,17401 | 0,0109284 | -0,002193 |
| $i=5$ | -0,14699 | -0,44704 | 0,8678151 | -0,22457 | 0,0115169 | -0,002980 |
| $i=6$ | -0,12573 | -0,55053 | 0,8730642 | -0,27059 | 0,0126608 | -0,003924 |
| $i=7$ | -0,08687 | -0,65307 | 0,8820957 | -0,30881 | 0,0145863 | -0,005106 |
| $i=8$ | -0,03235 | -0,75225 | 0,8936691 | -0,33773 | 0,0177924 | -0,006724 |
| $j=1$ | 0,066205 | 0,066896 | 0,9118716 | 0,027519 | 0,013097 | 0,000395 |
| $j=2$ | 0,139393 | 0,1265 | 0,9235324 | 0,048994 | 0,014282 | 0,000758 |
| $j=3$ | 0,217908 | 0,179553 | 0,934614 | 0,065433 | 0,0157261 | 0,001101 |
| $j=4$ | 0,300464 | 0,22683 | 0,9449365 | 0,077827 | 0,0174868 | 0,001440 |
| $j=5$ | 0,386083 | 0,269061 | 0,9544394 | 0,087037 | 0,0196698 | 0,001794 |
| $j=6$ | 0,474026 | 0,306907 | 0,9631315 | 0,093772 | 0,0224678 | 0,002187 |
| $j=7$ | 0,563739 | 0,34095 | 0,971058 | 0,098594 | 0,0262572 | 0,002666 |
| $j=8$ | 0,654805 | 0,371687 | 0,9782804 | 0,101939 | 0,0318835 | 0,003322 |
| String 3 |  |  |  |  | 0,2736473 | -0,010172 |
| Bodies 1 and 2 and string 3 |  |  |  |  | 9,8354767 | -0,700351 |

As a result of numerical calculation for the case, when the values of the coefficient of proportionality $\gamma_{V}\left(\right.$ also $\left.\gamma_{v}\right)$ lie in the range $\gamma_{V} \geq 1$ (and $\gamma_{v} \geq 1$ ),
it was found that, in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ at time $t_{1 p}$ the closed system of bodies 1 and 2 and string 3 has the projection of the momentum on the axis $\mathrm{O}_{1} x_{1}$, equal to $7,1214557 \cdot \boldsymbol{c}_{\boldsymbol{1}} \cdot \boldsymbol{M}_{\mathbf{0}}$, and the projection of the momentum on the axis $\mathrm{O}_{1} y_{1}$, equal to $\mathbf{0}$.

And in a stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ at time $t_{1 h}$ the closed system of bodies 1 and 2 and string 3 has the projection of the momentum on the axis $\mathrm{O}_{1} x_{1}$, equal to $\mathbf{9 , 8 3 5 4 7 6 7} \cdot \boldsymbol{c}_{\mathbf{1}} \cdot \boldsymbol{M}_{\mathbf{0}}$, and the projection of the momentum on the axis $\mathrm{O}_{1} y_{1}$, equal to $\mathbf{- 0 , 7 0 0 3 5 1} \cdot \boldsymbol{c}_{\mathbf{1}} \cdot \boldsymbol{M}_{\mathbf{0}}$.

As a result, we have a violation of the law of conservation of momentum for a closed mechanical system of bodies, because $7,1214557 \neq 9,8354767$ and $0 \neq-0,700351$.

Moreover, integration of mass of string 3 in calculating the momentum of the system of bodies 1 and 2 and string 3 leads to the aggravation of violating the law of conservation of momentum.

In the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ in the case, when the values of the coefficient of proportionality $\gamma_{V}$ (also $\gamma_{v}$ ) lie in the range $\gamma_{V} \geq 1$ (and $\gamma_{v} \geq 1$ ), the momentum of a closed system of bodies 1 and 2 and string 3 is not constant, as is a function of time $t_{1}$.

Next, consider the case, when the values of the coefficient of proportionality $\gamma_{V}$ (also $\gamma_{v}$ ) lie in the range $0<\gamma_{V} \leq 1$ (and $\left.0<\gamma_{v} \leq 1\right)$.

In the case, if the values of the coefficient of proportionality $\gamma_{v}$ are in the range $0<\gamma_{v} \leq 1$, then as follows from formula (75), in any inertial reference system $O x y z$ the projections $P_{x}$ and $P_{\mathrm{y}}$ of the momentum of a material point, moving with the speed $v$ and having a rest mass $m_{0}$, on the axis $O x$ and $O y$, respectively, can be written:

$$
\begin{equation*}
P_{x}=\frac{m_{\boldsymbol{o}} v_{x}}{\sqrt{1+\frac{\left(v_{x}^{2}+v_{y}^{2}\right)}{c_{2}^{2}}}} \tag{233}
\end{equation*}
$$

$$
P_{\mathrm{y}}=\frac{60}{\sqrt{\boldsymbol{o}} v_{y}} \sqrt{1+\frac{\left(v_{x}^{2}+v_{y}^{2}\right)}{c_{2}^{2}}}
$$

where: $v_{x}$ and $v_{y}$ - the projections of the speed $v$ of a material point on the axis $O x$ and $O y$, respectively.

Assume in the considered example 3 (shown in Fig. 13 - Fig.17), that:

$$
\begin{align*}
& \frac{V}{c_{2}}=0,9  \tag{235}\\
& \frac{v_{R}}{c_{2}}=0,8  \tag{236}\\
& \frac{m_{0}}{M_{0}}=0,1 \tag{237}
\end{align*}
$$

To determine the values of the momentums of the system of bodies 1 and 2 and string 3 in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ at time $t_{1 p}$ will use equation (99) - (102), (233)-(234), the raw data (214) - (216), (235)-(237) and the formulas, derived from the equations (174), (175), (181), (182), (189), (190), (197) and (198), taking into account equation (34) :

$$
\begin{align*}
v_{x 11}= & \frac{v_{x 21}+V}{1-\frac{V \cdot v_{x 21}}{c_{2}^{2}}}  \tag{238}\\
v_{y 11}= & \frac{v_{y 21} \cdot \sqrt{1+\frac{V^{2}}{c_{2}^{2}}}}{1-\frac{V \cdot v_{x 21}}{c_{2}^{2}}}  \tag{239}\\
v_{x 12}= & \frac{v_{x 22}+V}{1-\frac{V \cdot v_{x 22}}{c_{2}^{2}}}  \tag{240}\\
v_{y 12}= & \frac{1-\sqrt{1+\frac{V^{2}}{c_{2}^{2}}}}{v_{y 22}} \tag{241}
\end{align*}
$$

$$
\begin{align*}
& v_{x 11 \rho i}=\frac{61}{1-\frac{V \cdot v_{x 21 \rho i}}{c_{2}^{2}}} \\
& v_{y 11 \rho i}=\frac{v_{y 21 \rho i} \cdot \sqrt{1+\frac{V^{2}}{c_{2}^{2}}}}{1-\frac{V \cdot v_{x 21 \rho i}}{c_{2}^{2}}}  \tag{242}\\
& v_{x 12 \rho j}=\frac{v_{x 22 \rho j}+V}{1-\frac{V \cdot v_{x 22 \rho j}}{c_{2}^{2}}}  \tag{243}\\
& v_{y 22 \rho j} \cdot \sqrt{1+\frac{V^{2}}{c_{2}^{2}}}  \tag{244}\\
& v_{y 12 \rho j} \tag{245}
\end{align*}
$$

The results of digital calculations presented in Tab.9.

Range $0<\gamma_{V} \leq 1\left(\right.$ and $\left.0<\gamma_{v} \leq 1\right)$. Time $t_{1 p}$.

| Object | Mobile reference system <br> отсчета $\mathrm{O}_{2} x_{2} y_{2} z_{2}$$\|$The projections of the <br> velocity <br> (the dimension $c_{2}$ ) |  | Stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | The projections of the velocity <br> (the dimension $\mathrm{c}_{2}$ ) |  | $\begin{gathered} \text { Projections of the } \\ \text { momentum } \\ \text { (the dimension } \mathrm{c}_{2} \cdot M_{0} \text { ) } \end{gathered}$ |  |
|  | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{2} x_{2} \\ \hline \end{gathered}$ | on the axis $\mathrm{O}_{2} y_{2}$ | on the axis $\mathrm{O}_{1} x_{1}$ | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{1} y_{1} \end{gathered}$ | on the axis $\mathrm{O}_{1} x_{1}$ | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{1} y_{1} \end{gathered}$ |
| Body 1 | -0,8 | 0 | 0,0581395 | 0 | 0,0580415 | 0 |
| Body 2 | 0,8 | 0 | 6,0714286 | 0 | 0,9867059 | 0 |
| Bodies 1 and 2 |  |  |  |  | 1,0447474 | 0 |
| $i=0$ | 0 | 0 | 0,9 | 0 | 0,0039351 | 0 |
| $i=1$ | -0,09412 | 0 | 0,7429501 | 0 | 0,0035081 | 0 |
| $i=2$ | -0,18824 | 0 | 0,6086519 | 0 | 0,0030583 | 0 |
| $i=3$ | -0,28235 | 0 | 0,4924953 | 0 | 0,0025989 | 0 |
| $i=4$ | -0,37647 | 0 | 0,3910369 | 0 | 0,0021422 | 0 |
| $i=5$ | -0,47059 | 0 | 0,3016529 | 0 | 0,0016988 | 0 |
| $i=6$ | -0,56471 | 0 | 0,2223089 | 0 | 0,0012765 | 0 |
| $i=7$ | -0,65882 | 0 | 0,1514032 | 0 | 0,0008806 | 0 |
| $i=8$ | -0,75294 | 0 | 0,00876578 | 0 | 0,0005137 | 0 |
| $j=1$ | 0,094118 | 0 | 1,0861183 | 0 | 0,0043275 | 0 |
| $j=2$ | 0,188235 | 0 | 1,3101983 | 0 | 0,004676 | 0 |
| $j=3$ | 0,282353 | 0 | 1,5851735 | 0 | 0,0049751 | 0 |
| $j=4$ | 0,376471 | 0 | 1,930605 | 0 | 0,0052232 | 0 |
| $j=5$ | 0,470588 | 0 | 2,377551 | 0 | 0,0054223 | 0 |
| $j=6$ | 0,564706 | 0 | 2,9784689 | 0 | 0,0055764 | 0 |
| $j=7$ | 0,658824 | 0 | 3,8294798 | 0 | 0,0056915 | 0 |
| $j=8$ | 0,752941 | 0 | 5,1277372 | 0 | 0,0057736 | 0 |
| String 3 |  |  |  |  | 0,0612779 | 0 |
| Bodies 1 and 2 and string 3 |  |  |  |  | 1,106025 | 0 |

To determine the values of the momentums of the system of bodies 1 and 2
and string 3 in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ at time $t_{1 h}$ will use equation (155)-(156), (161)-(164), (111)-(112), (233)-(234), (238)-(245), the raw data (235)-(237), (214)-(216) and the formulas, derived from the equations (209), (211) and (213), taking into account equation (34) :

$$
\begin{array}{r}
\left(\omega \cdot t_{22 h}\right)=-\frac{V \cdot v_{R} \cdot\left[1+\cos \left(\omega \cdot t_{22 h}\right)\right]}{c_{2}^{2}} \\
\left(\omega \cdot t_{21 \rho i h}\right)=-\frac{V \cdot v_{R}}{c_{2}^{2}} \cdot\left\{1-\left[\frac{\rho_{i}}{R} \cdot \cos \left(\omega \cdot t_{21 \rho i h}\right)\right]\right\} \\
\left(\omega \cdot t_{22 \rho j h}\right)=-\frac{V \cdot v_{R}}{c_{2}^{2}} \cdot\left\{1+\left[\frac{\rho_{j}}{R} \cdot \cos \left(\omega \cdot t_{22 \rho j h}\right)\right]\right\} \tag{248}
\end{array}
$$

The results of digital calculations presented in Tab.10.

Tab. 10
Range $0<\gamma_{V} \leq 1\left(\right.$ and $\left.0<\gamma_{v} \leq 1\right)$. Time $t_{1 h}$.

| Object | Mobile reference system <br> $\mathrm{O}_{2} x_{2} y_{2} z_{2}$ |  | Stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | The projections of the velocity <br> (the dimension $\mathrm{c}_{2}$ ) |  | Projections of the momentum <br> (the dimension $\mathrm{c}_{2} \cdot M_{0}$ ) |  |
|  | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{2} x_{2} \end{gathered}$ | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{2} y_{2} \end{gathered}$ | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{1} x_{1} \end{gathered}$ | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{1} y_{1} \\ \hline \end{gathered}$ | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{1} x_{1} \end{gathered}$ | $\begin{gathered} \hline \text { on the axis } \\ \mathrm{O}_{1} y_{1} \\ \hline \end{gathered}$ |
| Body 1 | 0 | -0,8 | 0,9 | -1,07629 | 0,5223737 | -0,624695 |
| Body 2 | -0,700743 | 0,385953 | 0,1221935 | 0,318425 | 0,1156519 | 0,3013783 |
| Bodies 1 and 2 |  |  |  |  | 0,6380255 | -0,323317 |
| $i=0$ | 0 | 0 | 0,9 | 0 | 0,0039351 | 0 |
| $i=1$ | 0,05716 | -0,07477 | 1,0090732 | -0,10605 | 0,0041666 | -0,000438 |
| $i=2$ | 0,10286 | -0,15765 | 1,1051724 | -0,23373 | 0,0043091 | -0,000911 |
| $i=3$ | 0,13452 | -0,24825 | 1,177014 | -0,37999 | 0,004353 | -0,001405 |
| $i=4$ | 0,14977 | -0,3454 | 1,2133127 | -0,53708 | 0,0042956 | -0,001901 |
| $i=5$ | 0,14699 | -0,44704 | 1,2066039 | -0,69313 | 0,004142 | -0,002379 |
| $i=6$ | 0,12573 | -0,55053 | 1,1565986 | -0,83517 | 0,0039051 | -0,00282 |
| $i=7$ | 0,08687 | -0,65307 | 1,0705621 | -0,95313 | 0,0036032 | -0,003208 |
| $i=8$ | 0,03235 | -0,75225 | 0,9603103 | -1,04239 | 0,0032566 | -0,003535 |
| $j=1$ | -0,066205 | 0,066896 | 0,7869078 | 0,084938 | 0,0036296 | 0,000392 |
| $j=2$ | -0,139393 | 0,1265 | 0,6758229 | 0,151217 | 0,0032682 | 0,000731 |
| $j=3$ | -0,217908 | 0,179553 | 0,5702556 | 0,201957 | 0,0028701 | 0,001016 |
| $j=4$ | -0,300464 | 0,22683 | 0,4719207 | 0,240211 | 0,0024533 | 0,001249 |
| $j=5$ | -0,386083 | 0,269061 | 0,3813932 | 0,268639 | 0,0020331 | 0,001432 |
| $j=6$ | -0,474026 | 0,306907 | 0,2985891 | 0,289426 | 0,0016218 | 0,001572 |
| $j=7$ | -0,563739 | 0,34095 | 0,2230786 | 0,304306 | 0,0012277 | 0,001675 |
| $j=8$ | -0,654805 | 0,371687 | 0,1542765 | 0,314633 | 0,0008564 | 0,001747 |
| String 3 |  |  |  |  | 0,0539268 | -0,006784 |
| Bodies 1 and 2 and string 3 |  |  |  |  | 0,6919523 | -0,330101 |

As a result of numerical calculation for the case, when the values of the coefficient of proportionality $\gamma_{V}\left(\right.$ also $\left.\gamma_{v}\right)$ lie in the range $0<\gamma_{V}<1$ (and
$0<\gamma_{v}<1$ ), it was found, that in the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ at time $t_{1 p}$ the closed system of bodies 1 and 2 and string 3 has the projection of the momentum on the axis $\mathrm{O}_{1} x_{1}$, equal to $\mathbf{1 , 1 0 6 0 2 5} \cdot \boldsymbol{c}_{2} \cdot \boldsymbol{M}_{\mathbf{0}}$, and the projection of the momentum on the axis $\mathrm{O}_{1} y_{1}$, equal to $\mathbf{0}$.

And in a stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ at time $t_{1 h}$ the closed system of bodies 1 and 2 and string 3 has the projection of the momentum on the axis $\mathrm{O}_{1} x_{1}$, equal to $\mathbf{0 , 6 9 1 9 5 2 3} \cdot \boldsymbol{c}_{\mathbf{2}} \cdot \boldsymbol{M}_{\mathbf{0}}$, and the projection of the momentum on the axis $\mathrm{O}_{1} y_{1}$, equal to $\mathbf{- 0 , 3 3 0 1 0 1} \cdot \boldsymbol{c}_{\mathbf{2}} \cdot \boldsymbol{M}_{\mathbf{0}}$.

As a result, we have a violation of the law of conservation of momentum for a closed mechanical system of bodies, because $1,106025 \neq 0,6919523$ and $0 \neq-0,330101$.

Moreover, integration of mass of string 3 in calculating the momentum of the system of bodies 1 and 2 and string 3 also leads to the aggravation of violating the law of conservation of momentum.

In the stationary reference system $\mathrm{O}_{1} x_{1} y_{1} z_{1}$ in the case, when the values of the coefficient of proportionality $\gamma_{V}$ (also $\gamma_{v}$ ) lie in the range $0<\gamma_{V}<1$ (and $0<\gamma_{v}<1$ ), the momentum of a closed system of bodies 1 and 2 and string 3 is also not constant, as is a function of time $t_{1}$.

## 9. Conclusion

In conclusion, we note the following:

- there are possible two variants of the relationship between coordinates and time in inertial reference systems for values of the coefficient of proportionality $\gamma_{V}$, are in the range $\gamma_{V}>1$ and $0<\gamma_{V}<1$.
- use of the special theory of relativity in dealing with individual examples (examples 2 and 3 ) may lead to non-compliance with the law of conservation of momentum for a closed mechanical system in the inertial reference systems.

Given, that the law of conservation of momentum associated with the homogeneity of space, we can assume, that the failure of the law of conservation of momentum will lead to non-compliance with conditions of symmetry of
space and time, on which is based the special theory of relativity.
The results obtained, when considering the examples 2 and 3 show that, if true to the law of conservation of momentum, it is necessary to finalize the special theory of relativity, or if it is true the special theory of relativity, then, consequently, incorrect law of conservation of momentum - possible to change the momentum of a closed system over time in the inertial reference systems.

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